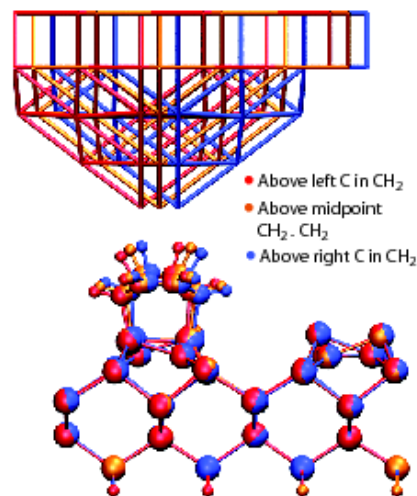
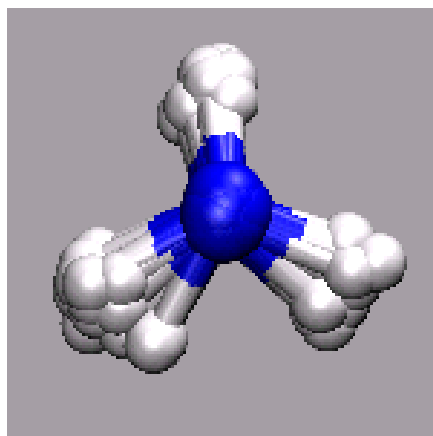
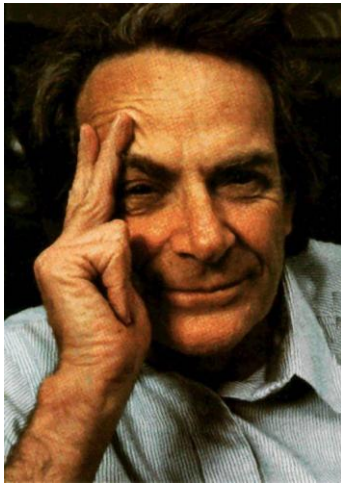


Two (unrelated) topics in quantum transport: Evaluating centroid and ring-polymer molecular dynamics and calculating STM images for molecules on semiconductor surfaces



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Quantum Statistical Mechanics á la Feynman

Energy eigenvalues:

$$\hat{H} |E_k\rangle = E_k |E_k\rangle$$

Canonical partition function:

$$Q = \text{Tr} \left[e^{-\beta \hat{H}} \right] = \sum_k e^{-\beta E_k}$$

Ensemble average:

$$\langle \hat{A} \rangle = \frac{1}{Q} \text{Tr} \left[\hat{A} e^{-\beta \hat{H}} \right] = \frac{1}{Q} \sum_k e^{-\beta E_k} \langle E_k | \hat{A} | E_k \rangle$$

“This law is the summit of statistical mechanics, and the entire subject is either the slide-down from this summit, as the principle is applied to various cases, or the climb-up to where the fundamental law is derived and the concepts of thermal equilibrium and temperature T clarified”

-- *Statistical Mechanics: A Set of Lectures*

Path Integrals in Quantum Statistical Mechanics

Carry out trace in coordinate basis:

$$Q = \text{Tr} \left[e^{-\beta \hat{H}} \right] = \int dx \langle x | e^{-\beta(\hat{K} + \hat{V})} | x \rangle$$

Trotter theorem:

$$e^{-\beta(\hat{K} + \hat{V})} = \lim_{P \rightarrow \infty} \left[e^{-\beta \hat{V} / P} e^{-\beta \hat{K} / P} \right]^P \equiv \lim_{P \rightarrow \infty} \hat{\Omega}^P$$

Insert completeness of position eigenvectors:

$$\begin{aligned} Q &= \lim_{P \rightarrow \infty} \int dx_1 \langle x_1 | \hat{\Omega}^P | x_1 \rangle \\ &= \lim_{P \rightarrow \infty} \int dx_1 dx_2 dx_3 \cdots dx_P \langle x_1 | \hat{\Omega} | x_2 \rangle \langle x_2 | \hat{\Omega} | x_3 \rangle \langle x_3 | \hat{\Omega} \cdots | x_P \rangle \langle x_P | \hat{\Omega} | x_1 \rangle \end{aligned}$$

Matrix elements:

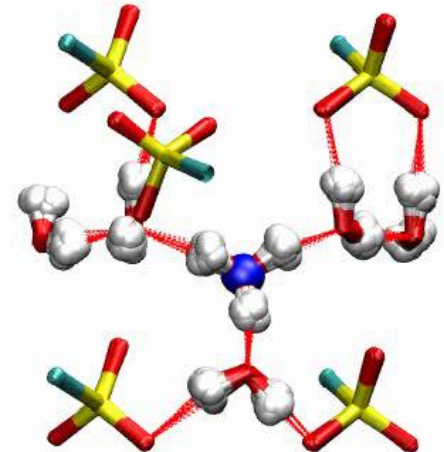
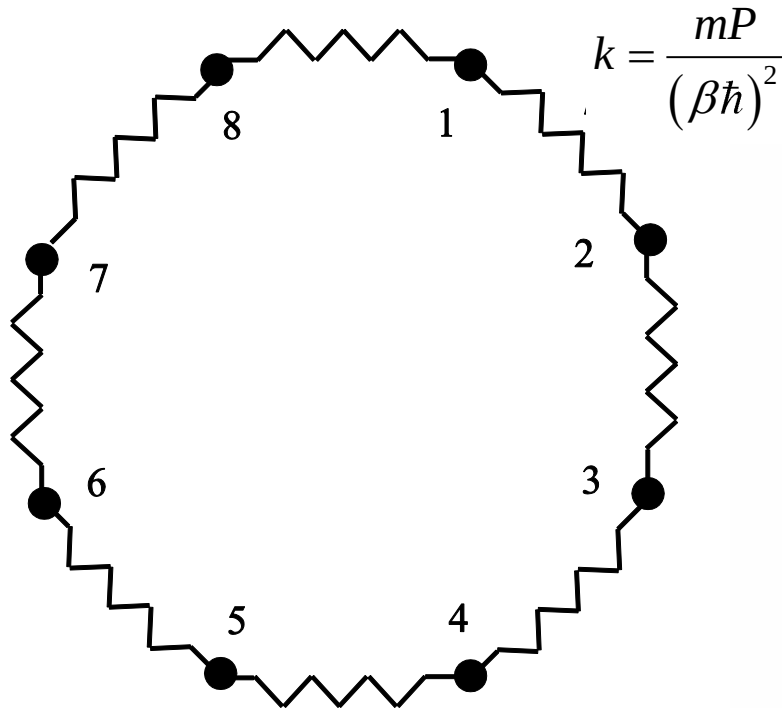
$$\langle x_i | \hat{\Omega} | x_{i+1} \rangle = \left(\frac{mP}{2\pi\beta\hbar^2} \right)^{1/2} \exp \left[-\frac{mP}{2\beta\hbar^2} (x_i - x_{i+1})^2 - \frac{\beta}{P} V(x_i) \right]$$

Discrete Path Integral

$$Q = \lim_{P \rightarrow \infty} Q_P$$

$$Q_P = \left(\frac{mP}{2\pi\beta\hbar^2} \right)^{P/2} \int dx_1 \cdots dx_P \exp \left\{ -\beta \sum_{i=1}^P \left[\frac{mP}{2\beta^2\hbar^2} (x_i - x_{i+1})^2 + \frac{1}{P} V(x_i) \right] \right\} \Bigg|_{x_{P+1}=x_1}$$

Classical
isomorphism:



Path integral molecular dynamics

M. Parrinello and A. Rahman, *J. Chem. Phys.* **80**, 860 (1984); R. W. Hall and B. J. Berne, *J. Chem. Phys.* **81**, 3641 (1984).

Add Gaussian integrals to partition function:

$$Q_P = \int dp_1 \cdots dp_P dx_1 \cdots dx_P \exp \left\{ -\beta \sum_{i=1}^P \left[\frac{p_i^2}{2m'} + \frac{1}{2} m \omega_P^2 (x_i - x_{i+1})^2 + \frac{1}{P} V(x_i) \right] \right\} \bigg|_{x_{P+1}=x_1}$$
$$\omega_P = \frac{\sqrt{P}}{\beta \hbar}$$

Partition function can be sampled using the Hamiltonian:

$$H = \sum_{i=1}^P \left[\frac{p_i^2}{2m'} + \frac{1}{2} m \omega_P^2 (x_i - x_{i+1})^2 + \frac{1}{P} V(x_i) \right]_{x_{P+1}=x_1}$$

“Primitive” equations of motion:

$$\dot{x}_i = \frac{p_i}{m'}, \quad \dot{p}_i = -m \omega_P^2 (2x_i - x_{i+1} - x_{i-1}) - \frac{1}{P} \frac{\partial V}{\partial x_i} + \text{Heat bath}$$

Variable Transformations in path integrals

MET, G. J. Martyna, M. L. Klein and B. J. Berne *J. Chem. Phys.* **99**, 2796 (1993)

Staging transformation:

$$u_1 = x_1$$

$$u_k = x_k - \frac{(k-1)x_{k+1} + x_1}{k}, \quad k = 2, \dots, P$$

Normal mode transformation:

$$x_k = \sum_{l=1}^P a_l e^{2\pi i(k-1)(l-1)/P}$$

$$u_1 = a_1, \quad u_P = a_{(P+2)/2}, \quad u_{2k-2} = \operatorname{Re}(a_k), \quad u_{2k-1} = \operatorname{Im}(a_k)$$

Diagonalization of quantum kinetic energy:

$$\sum_{k=1}^P (x_k - x_{k+1})^2 = \sum_{k=2}^P \lambda_k u_k^2$$

$$\text{Staging: } \lambda_k = \frac{k}{k-1} \quad \text{Normal modes: } \lambda_{2k-1} = \lambda_{2k-2} = 2 \left[1 - \cos \left(\frac{2\pi(k-1)}{P} \right) \right]$$

Path integral molecular dynamics

MET, G. J. Martyna, M. L. Klein and B. J. Berne *J. Chem. Phys.* **99**, 2796 (1993)

Transformed Hamiltonian:

$$H = \sum_{i=1}^P \left[\frac{p_i^2}{2m'_i} + \frac{1}{2} m \omega_P^2 \lambda_i u_i^2 + \frac{1}{P} V(x_i(u)) \right]_{x_{P+1}=x_1}$$

Fictitious masses:

$$m'_1 = m, \quad m'_i = \lambda_i m, \quad i = 2, \dots, P$$

Equations of motion:

$$\dot{u}_i = \frac{p_i}{m'_i}$$

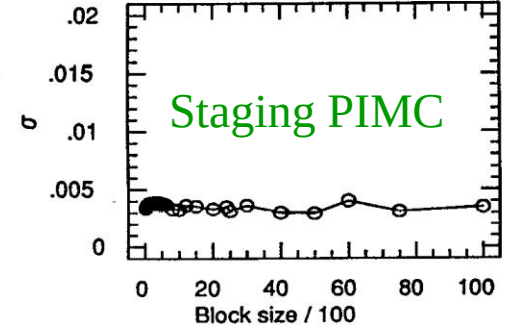
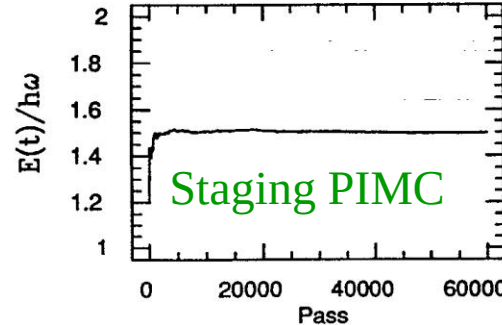
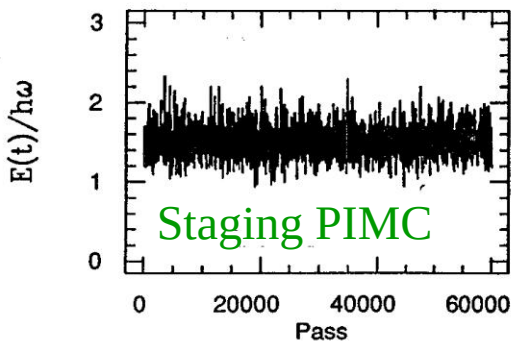
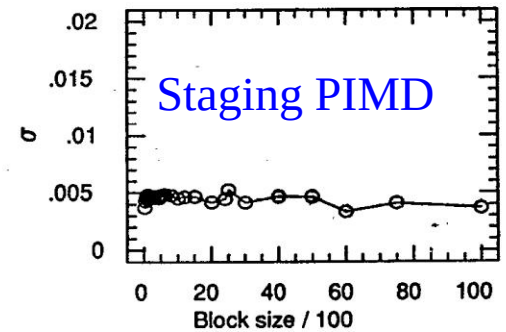
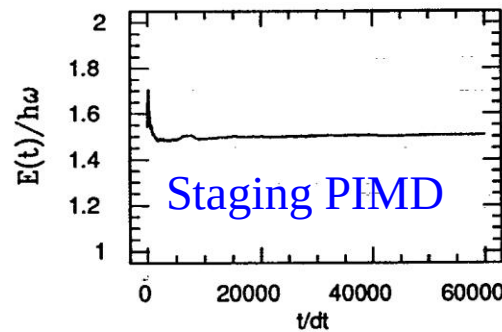
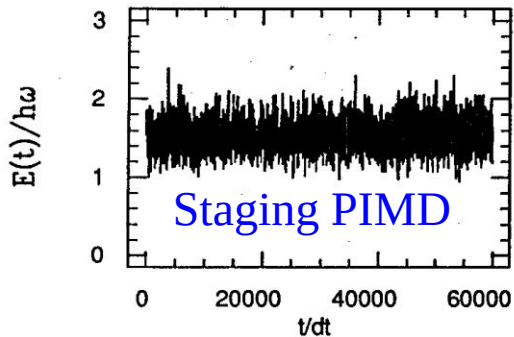
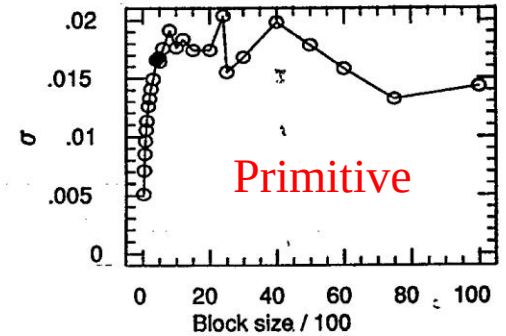
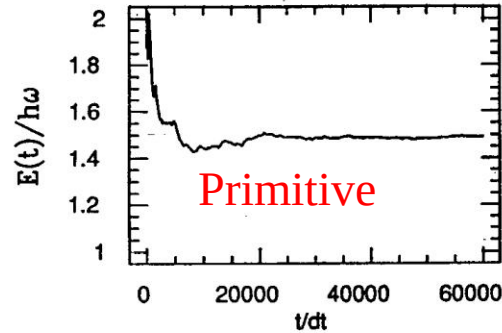
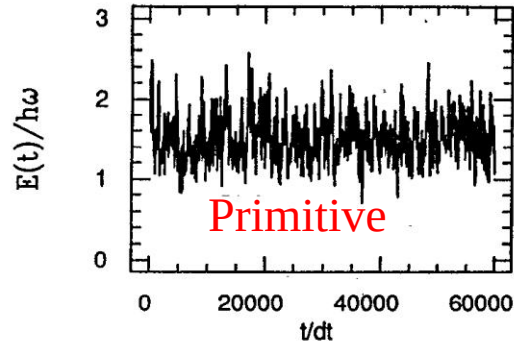
$$\dot{p}_i = -m \omega_P^2 \lambda_i u_i - \frac{1}{P} \sum_{j=1}^P \frac{\partial V}{\partial x_j} \frac{\partial x_j}{\partial u_i} + \text{Heat bath}(i)$$

Results for harmonic oscillator

$$H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2, \quad \beta\hbar\omega = 15.8, \quad \frac{m\omega}{\hbar} = 0.03, \quad P = 400$$

$$E_{\text{vir}} = \frac{1}{P} \sum_{i=1}^P \left[V(x_i) + \frac{1}{2} x_i \frac{\partial V}{\partial x_i} \right]$$

M. F. Herman, E. J. Bruskin and B. J. Berne
J. Chem. Phys. **76**, 5150 (1982).



A. Pérez, M. Müser, MET
J. Chem. Phys. **130**, 184105 (2009)

$$V(x) = \frac{x^4}{4}$$

$$\beta = 8, \quad P = 32$$

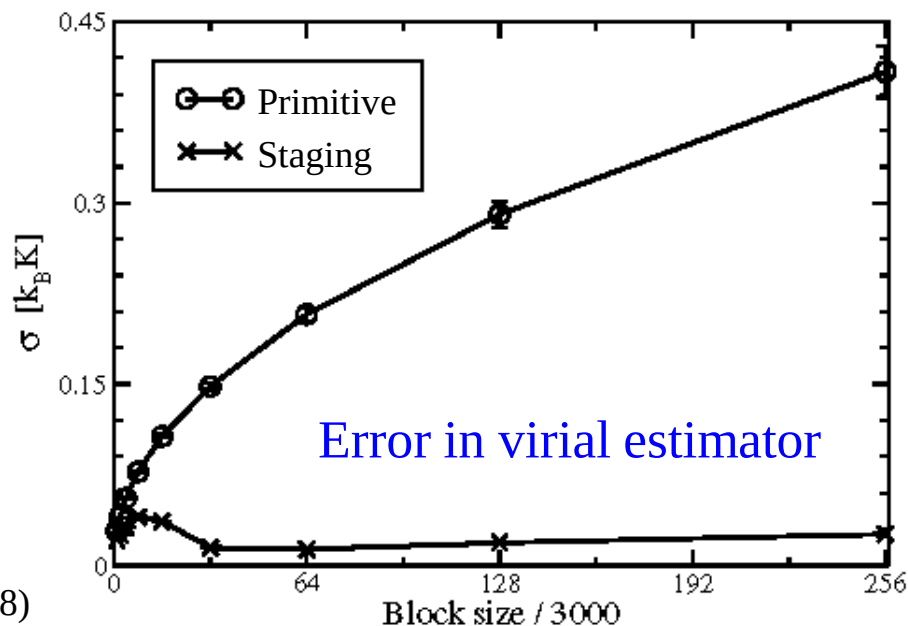
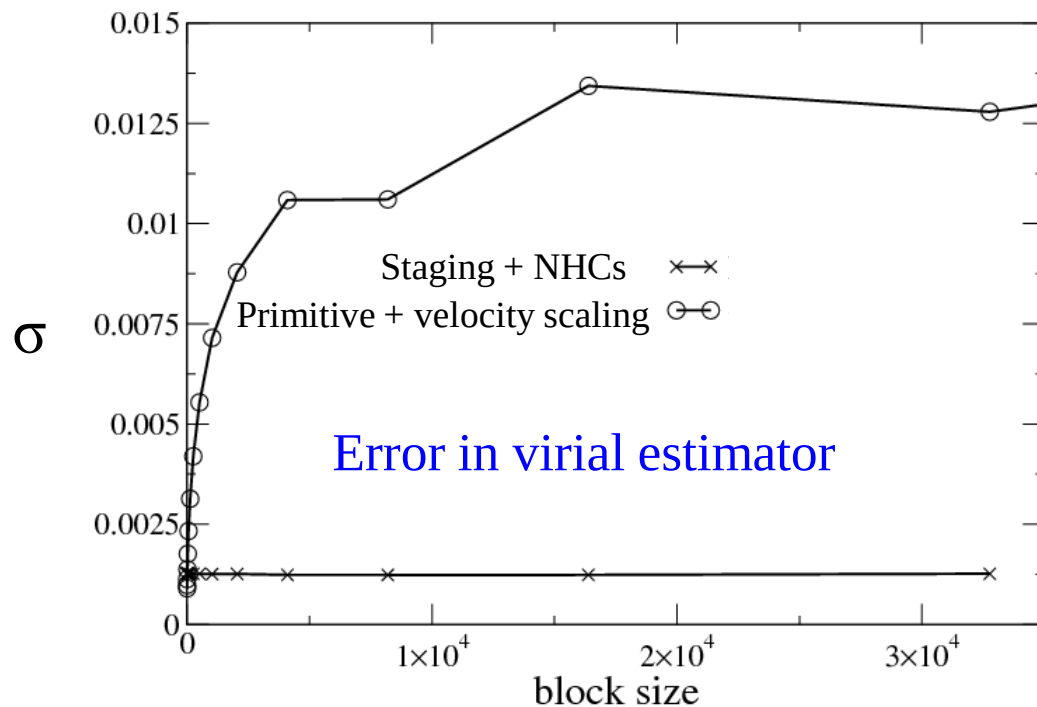
Liquid *para*-Hydrogen
 Silvera-Goldman potential

$$T = 14 \text{ K}, \quad \rho = 0.0235 \text{ \AA}^{-3}$$

$$N = 256, \quad P = 64$$

Time = 6 ps

I. F. Silvera and V. V. Goldman, *JCP* **69**, 4209 (1978)



Quantum time correlation functions

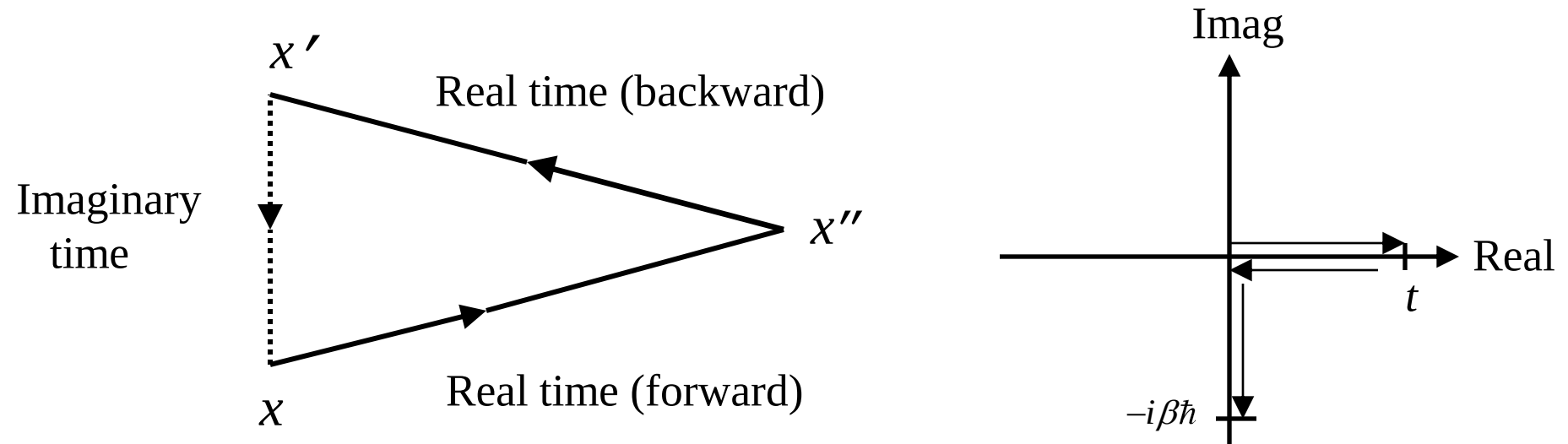
G. Krilov, E. Sim and B. J. Berne *J. Chem. Phys.* **114**, 1075 (2001)

Non-symmetric correlation function:

$$C_{AB}(t) = \langle \hat{A}(0) \hat{B}(t) \rangle = \frac{1}{Q} \text{Tr} \left[e^{-\beta \hat{H}} \hat{A} e^{i\hat{H}t/\hbar} \hat{B} e^{-i\hat{H}t/\hbar} \right]$$

Assume $\hat{A} = \hat{A}(\hat{x})$, $\hat{B} = \hat{B}(\hat{x})$, trace in coordinate basis

$$C_{AB}(t) = \frac{1}{Q} \int dx \, dx' \, dx'' \langle x | e^{-\beta \hat{H}} | x' \rangle a(x') \langle x' | e^{i\hat{H}t/\hbar} | x'' \rangle b(x'') \langle x'' | e^{-i\hat{H}t/\hbar} | x \rangle$$



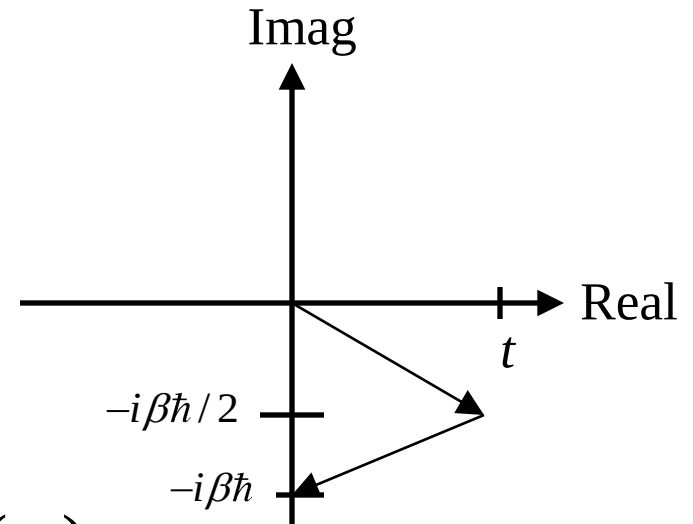
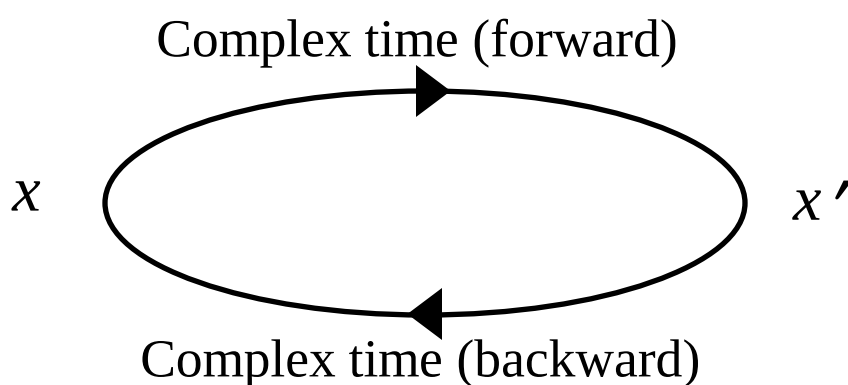
Quantum time correlation functions

Symmetric complex-time correlation function:

$$G_{AB}(t) = \frac{1}{Q} \text{Tr} \left[A e^{i\hat{H}\tau^*/\hbar} \hat{B} e^{-i\hat{H}\tau/\hbar} \right], \quad \tau = t - \frac{i\beta\hbar}{2}$$

Assume $\hat{A} = \hat{A}(\hat{x})$, $\hat{B} = \hat{B}(\hat{x})$, trace in coordinate basis

$$G_{AB}(t) = \frac{1}{Q} \int dx dx' a(x) \langle x | e^{i\hat{H}\tau^*/\hbar} | x' \rangle b(x') \langle x' | e^{-i\hat{H}\tau/\hbar} | x \rangle$$



Fourier transforms: $\tilde{C}_{AB}(\omega) = e^{\beta\hbar\omega} \tilde{G}_{AB}(\omega)$

Quantum time correlation functions

Kubo-transformed correlation function:

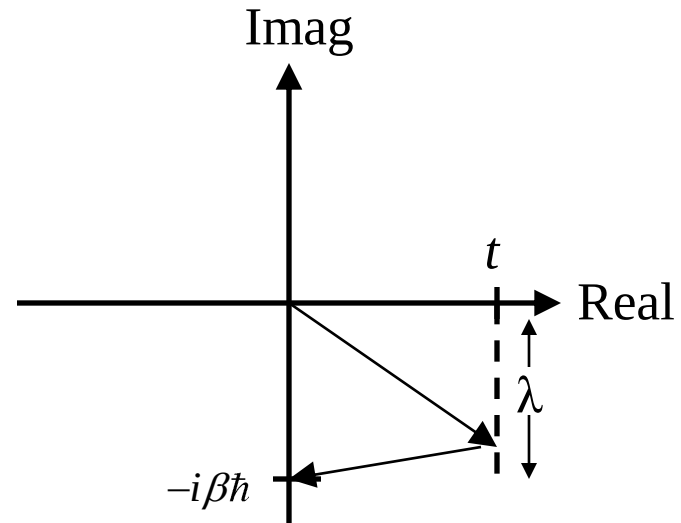
$$K_{AB}(t) = \frac{1}{\beta Q} \int_0^\beta d\lambda \operatorname{Tr} \left[A e^{i\hat{H}\tau_\lambda^*/\hbar} \hat{B} e^{-i\hat{H}\tau_{\beta-\lambda}/\hbar} \right], \quad \tau_\lambda = t - i\lambda\hbar$$

Assume $\hat{A} = \hat{A}(\hat{x})$, $\hat{B} = \hat{B}(\hat{x})$, trace in coordinate basis

$$K_{AB}(t) = \frac{1}{\beta Q} \int_0^\beta d\lambda \int dx dx' a(x) \langle x | e^{i\hat{H}\tau_\lambda^*/\hbar} | x' \rangle b(x') \langle x' | e^{-i\hat{H}\tau_{\beta-\lambda}/\hbar} | x \rangle$$

Fourier transforms:

$$\tilde{C}_{AB}(\omega) = \left[\frac{\beta\hbar\omega}{1 - e^{-\beta\hbar\omega}} \right] \tilde{K}_{AB}(\omega)$$



Centroid molecular dynamics

J. Cao and G. A. Voth *J. Chem. Phys.* **99**, 10070 (1993)

J. Cao and G. A. Voth *J. Chem. Phys.* **100**, 5106 (1994)

J. Cao and G. A. Voth *J. Chem. Phys.* **101**, 6168 (1994)

Path Centroid

$$x_c = \frac{1}{P} \sum_{i=1}^P x_i$$

x_c is the unbound mode u_1 in a normal-mode transformation.

Centroid potential of mean force: (Feynman and Kleinert *Phys. Rev. A* **34**, 5080 (1986))

$$W(x_c) = W(u_1) = -kT \ln \left\{ \int du_2 \cdots du_P \exp \left[-\beta \sum_{k=2}^P \frac{1}{2} m \omega_P^2 \lambda_k u_k^2 - \frac{\beta}{P} \sum_{k=1}^P V(x_k(u)) \right] \right\}$$

Discrete partitionfunction:

$$Q_P = \left(\frac{mP}{2\pi\beta\hbar^2} \right)^{P/2} \int dx_c e^{-\beta W(x_c)}$$

Centroid molecular dynamics

Approximate quantum dynamics as dynamics on the centroid PMF

$$\dot{x}_c = \frac{p_c}{m}, \quad \dot{p}_c = -\frac{dW}{dx_c}$$

Approximation to Kubo-transformed time correlation function:

$$K_{AB}(t) \approx \frac{1}{Q} \int dp_c dx_c a(x_c(0)) b(x_c(t)) \exp \left[-\beta \left(\frac{p_c^2}{2m} + W(x_c) \right) \right]$$

Generate centroid PMF “on the fly” by performing imaginary-time PIMD with

$$m'_1 = m, \quad m'_k = \gamma m_k, \quad \gamma < 1$$

called “partially adiabatic CMD”. (J. Cao and G. J. Martyna, *J. Chem. Phys.* **104**, 2028 (1996))

Short-time accuracy (B. J. Braams and D. E. Manolopoulos *J. Chem. Phys.* **127**, 174108 (2007))

$$K_{xx}(t) = \langle x_c(0)x_c(t) \rangle + O(t^6), \quad K_{vv}(t) = \langle v_c(0)v_c(t) \rangle + O(t^4)$$

Exact in harmonic and classical limits.

Ring-polymer molecular dynamics

I. R. Craig and D. E. Manolopoulos *J. Chem. Phys.* **121**, 3368 (2004)

Classical ring-polymer Hamiltonian:

$$H_{\text{cl}}(x, p) = \sum_{i=1}^P \left[\frac{p_i^2}{2m} + \frac{1}{2} m \omega_P^2 (x_i - x_{i+1})^2 + V(x_i) \right]_{x_{P+1}=x_1}$$

(Can also be written in terms of normal-mode and staging variables)

Equations of motion:

$$\dot{x}_i = \frac{p_i}{m}, \quad \dot{p}_i = -m \omega_P^2 (2x_i - x_{i+1} - x_{i-1}) - \frac{\partial V}{\partial x_i}$$

Approximate correlation function:

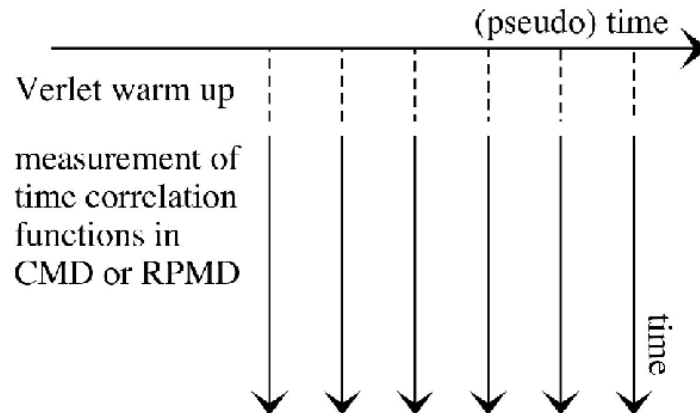
$$K_{AB}(t) \approx \frac{1}{Q} \int d^P p d^P x a_P(0) b_P(t) \exp[-\beta H_{\text{cl}}(x, p) / P]$$
$$a_P(t) = \frac{1}{P} \sum_{i=1}^P a(x_i(t)), \quad b_P(t) = \frac{1}{P} \sum_{i=1}^P b(x_i(t))$$

Short-time accuracy: (B. J. Braams and D. E. Manolopoulos *J. Chem. Phys.* **127**, 174108 (2007))

$$K_{xx}(t) = \langle x_P(0) x_P(t) \rangle + O(t^8), \quad K_{vv}(t) = \langle v_P(0) v_P(t) \rangle + O(t^6)$$

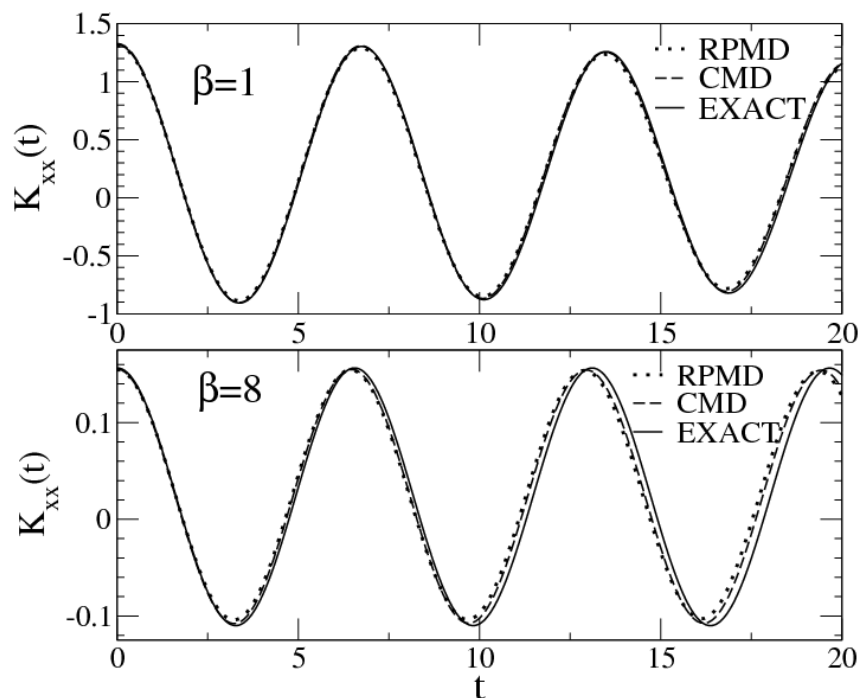
Some numerical examples

*Recommended
algorithm:*



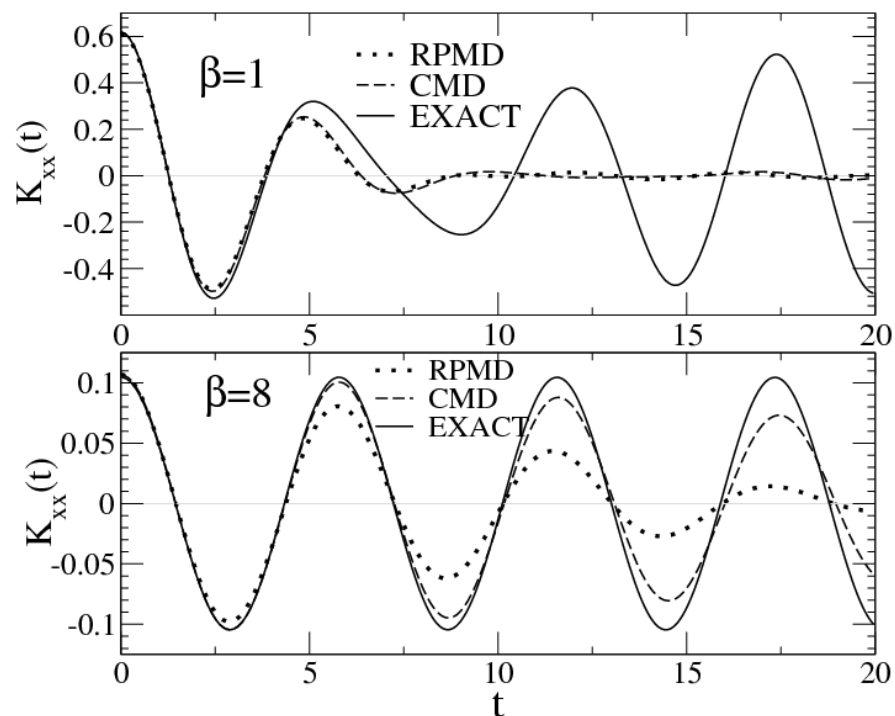
$$V(x) = \frac{x^2}{2} + 0.1x^3 + 0.01x^4$$

$$P = 32, \quad \gamma_{\text{CMD}} = 0.005$$

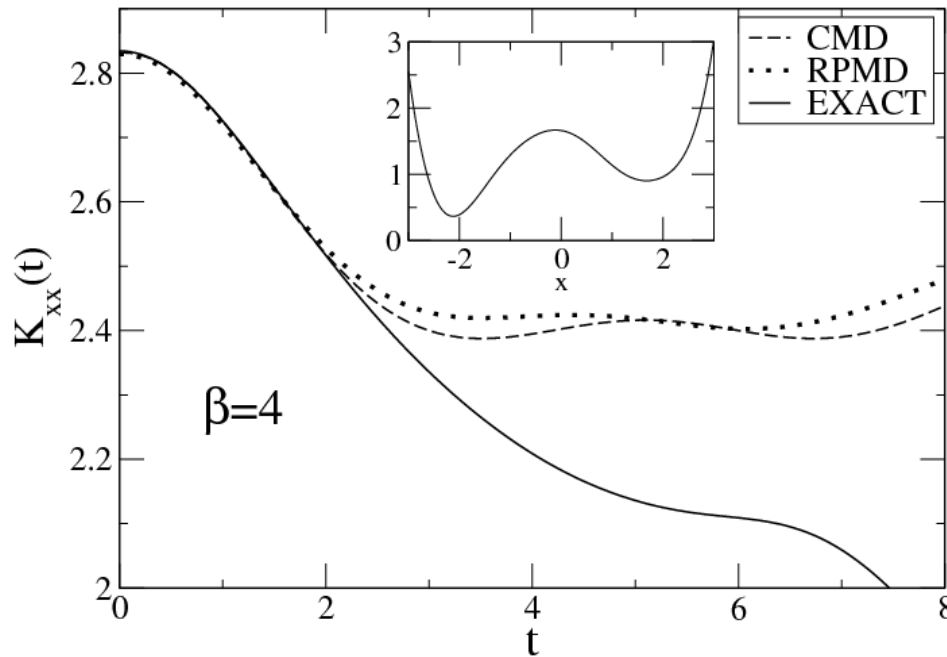


$$V(x) = \frac{x^4}{4}$$

$$P = 32, \quad \gamma_{\text{CMD}} = 0.005$$



Some numerical examples



$\beta=4$

$$V(x) = 0.1(x^2 - 4)^2 + e^{-(x-2.2)^2} + 0.5e^{-(x+1.5)^2}$$

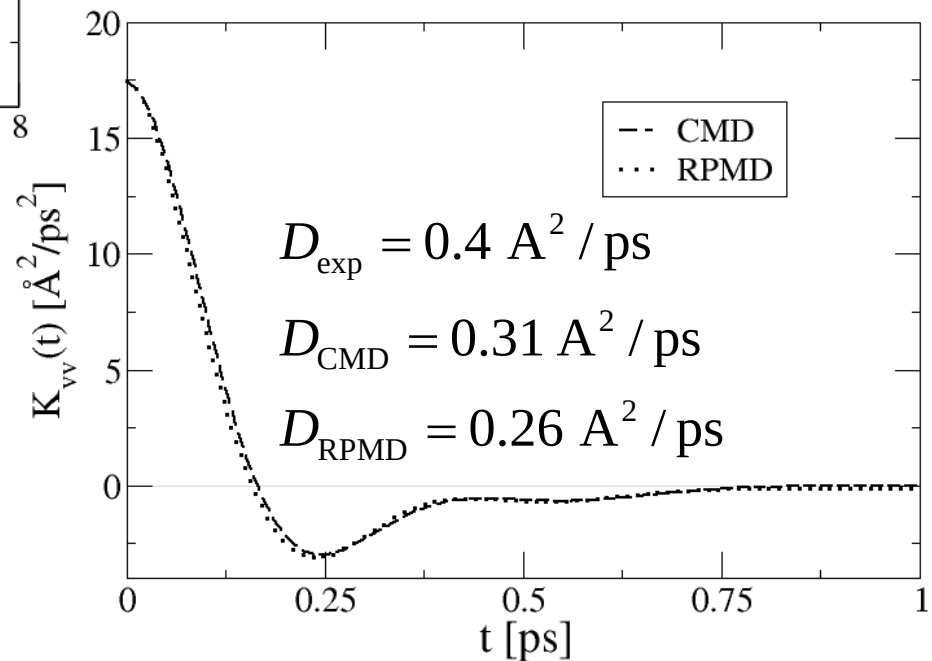
$$P = 32, \quad \gamma_{\text{CMD}} = 0.005$$

Liquid *para*-Hydrogen

$$T = 14 \text{ K}, \quad \rho = 0.0235 \text{ \AA}^{-3}$$

$$N = 256, \quad P = 64, \quad \gamma_{\text{CMD}} = 0.005$$

Time = 6 ps



$$D_{\text{exp}} = 0.4 \text{ \AA}^2 / \text{ps}$$

$$D_{\text{CMD}} = 0.31 \text{ \AA}^2 / \text{ps}$$

$$D_{\text{RPMD}} = 0.26 \text{ \AA}^2 / \text{ps}$$

Self-consistent quality measure of CMD and RPMD correlation functions

Imaginary-time mean-square displacement:

$$G(\tau) = \left\langle \left(\hat{x}(0) - \hat{x}(\tau) \right)^2 \right\rangle \xrightarrow{\text{free particle}} \frac{1}{2m\beta} \tau(\beta\hbar - \tau)$$

$$G(\tau) = \frac{1}{N} \sum_{i=1}^N \left\langle \left(\hat{\mathbf{r}}_i(0) - \hat{\mathbf{r}}_i(\tau) \right)^2 \right\rangle \xrightarrow{\text{free particles}} \frac{\pi}{m\beta} \tau(\beta\hbar - \tau)$$

Relation to real-time position and velocity correlation functions:

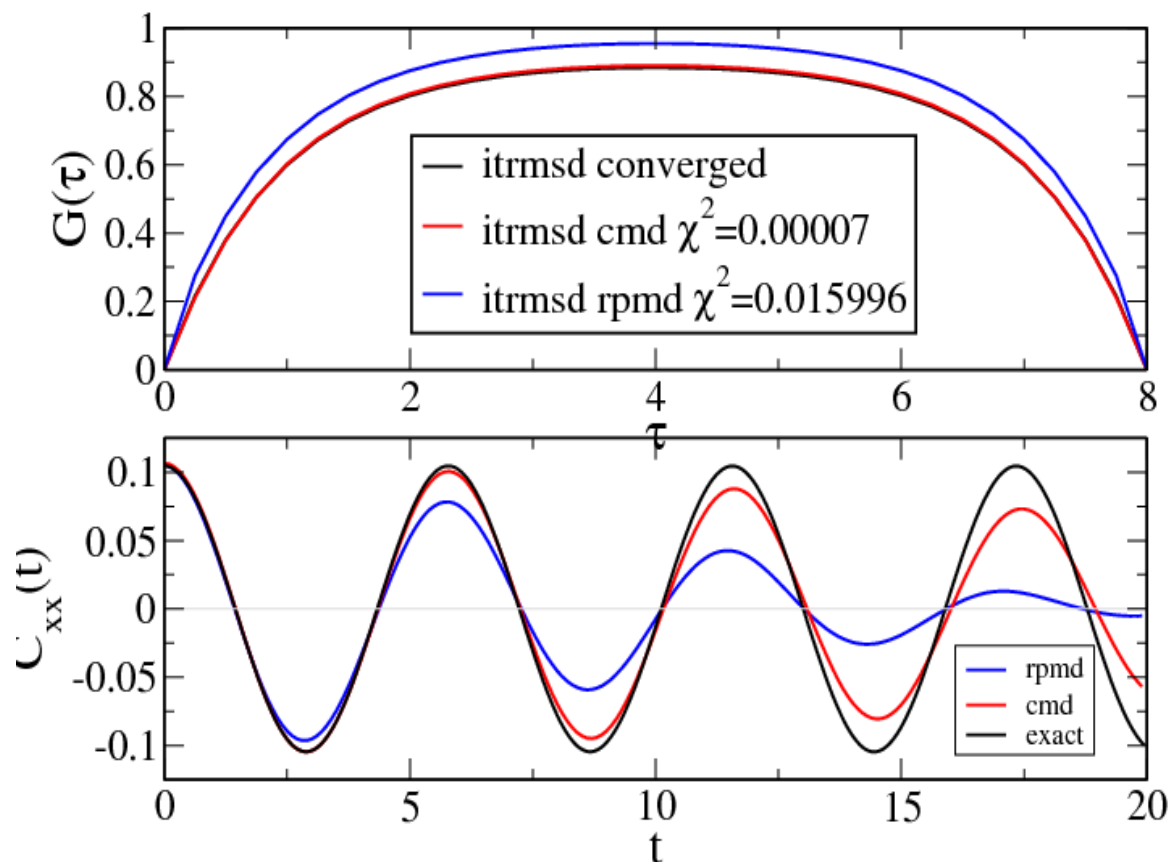
$$\begin{aligned} G(\tau) &= \frac{1}{\pi} \int_{-\infty}^{\infty} d\omega e^{-\beta\hbar\omega} \tilde{K}_{xx}(\omega) \left\{ \cosh \left[\omega \left(\frac{\beta\hbar}{2} - \tau \right) \right] - \cosh \left(\frac{\beta\hbar\omega}{2} \right) \right\} \\ &= \frac{1}{\pi} \int_{-\infty}^{\infty} d\omega e^{-\beta\hbar\omega} \frac{\tilde{K}_{vv}(\omega)}{\omega^2} \left\{ \cosh \left[\omega \left(\frac{\beta\hbar}{2} - \tau \right) \right] - \cosh \left(\frac{\beta\hbar\omega}{2} \right) \right\} \end{aligned}$$

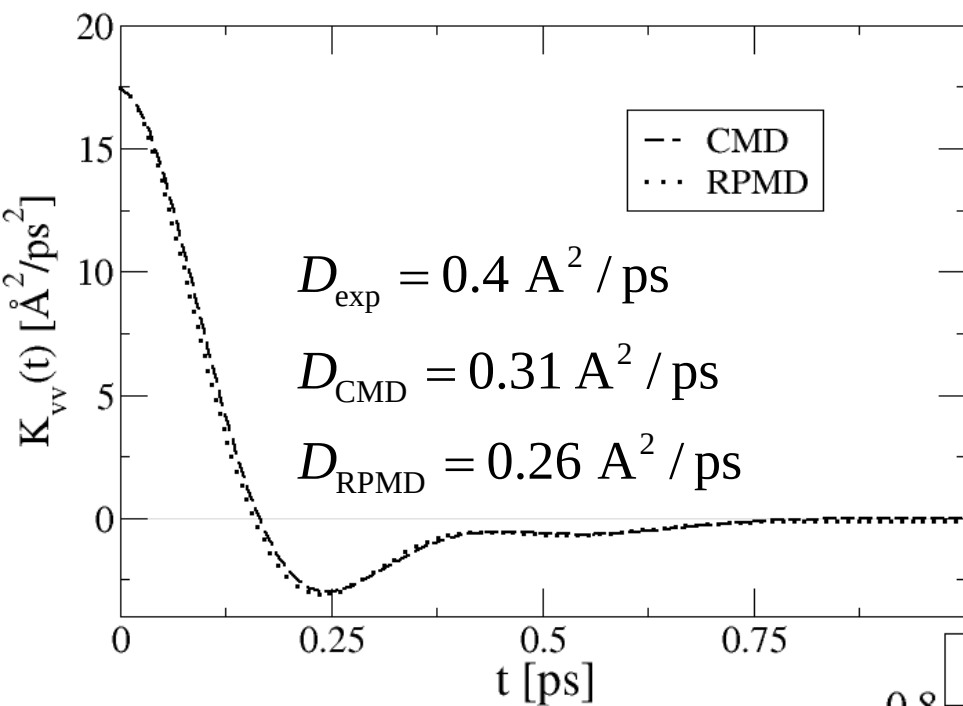
χ^2 Error Measure

A. Pérez, M. Müser, MET *J. Chem. Phys.* **130**, 184105 (2009)

$$\chi^2 = \frac{1}{\beta\hbar} \int_0^{\beta\hbar} d\tau \left[\frac{G^{(\text{est})}(\tau) - G(\tau)}{G(\tau)} \right]^2$$

ITRMSD of $0.25X^4$ with $P=32$ and $\beta=8$





Liquid *para*-Hydrogen

$T = 14 \text{ K}$, $\rho = 0.0235 \text{ \AA}^{-3}$

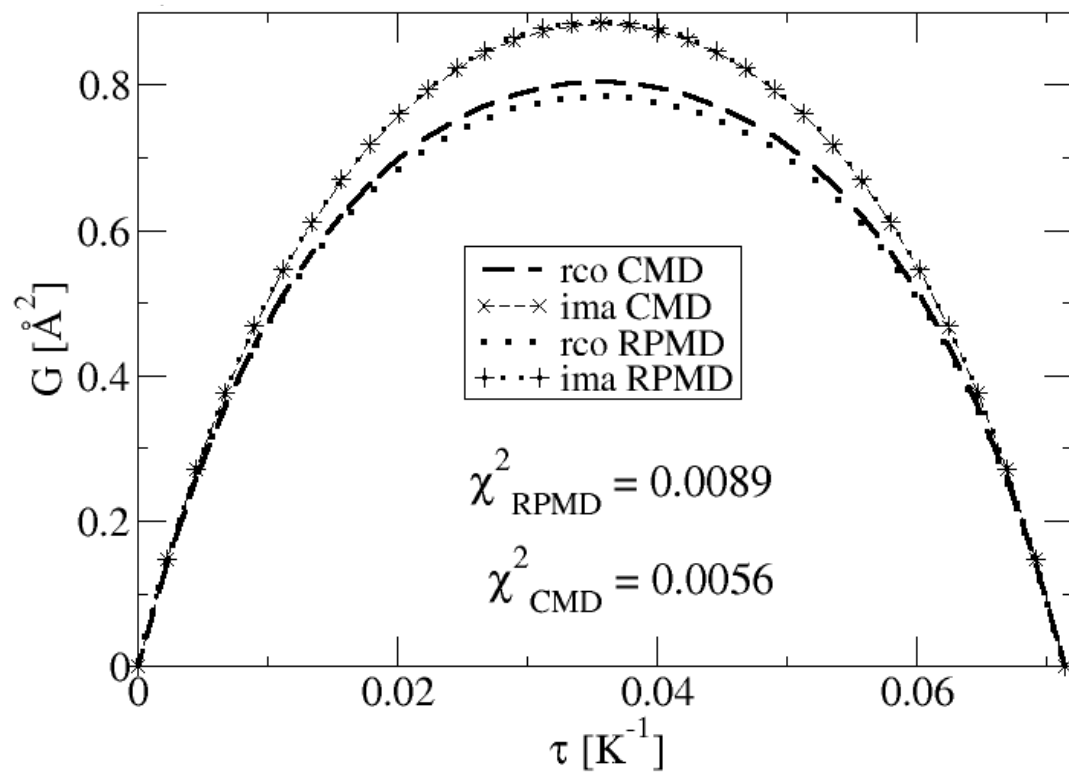
$N = 256$, $P = 64$, $\gamma_{\text{CMD}} = 0.005$

Time = 6 ps

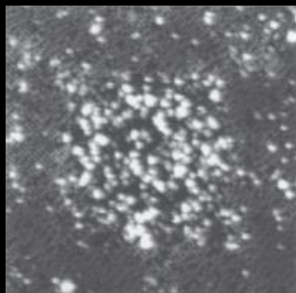
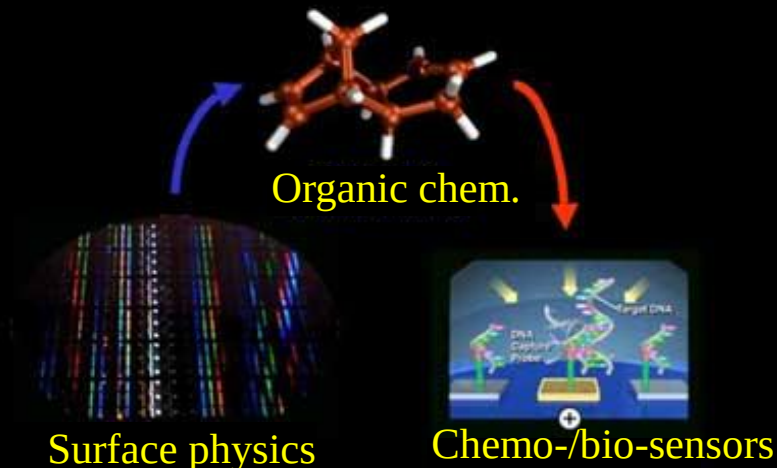
Considerations of CMD and RPMD
in calculating vibrational spectra:

A. Witt *et al.* *JCP* **130**, 194501 (2009)

S. D. Ivanov *et al.* *JCP* **132**, 031101 (2010)



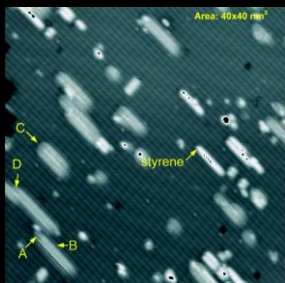
Organic-Semiconductor Interfaces



Nano-lithography

- Example:
- Passivate Si(100) surface with benzene
 - Create 2 nm wide patterns with STM tip
 - React with vinyl ferrocene

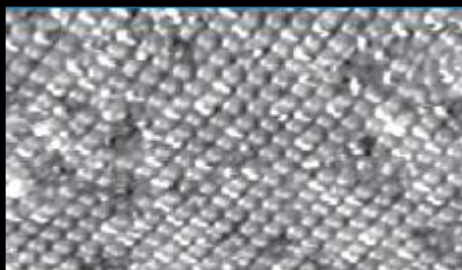
Kruse and Wolkow *Appl. Phys. Lett.* **81** (2002) 4422.



Self-assembled nanowires and other nanostructures

- Example:
- styrene forms lines on H-Si(100)
 - precursor to molecular electronics

DiLabio, Piva, Kruse, and Wolkow *JACS* **126** (2004) 16048.



Monolayers

- Example:
- monolayer of 1,5-cyclooctadiene absorbed on Si(100)
 - π -bond on surface available for further rxns
 - precursor to molecular sensor

DiLabio, Piva, Kruse, and Wolkow *JACS* **126** (2004) 16048.

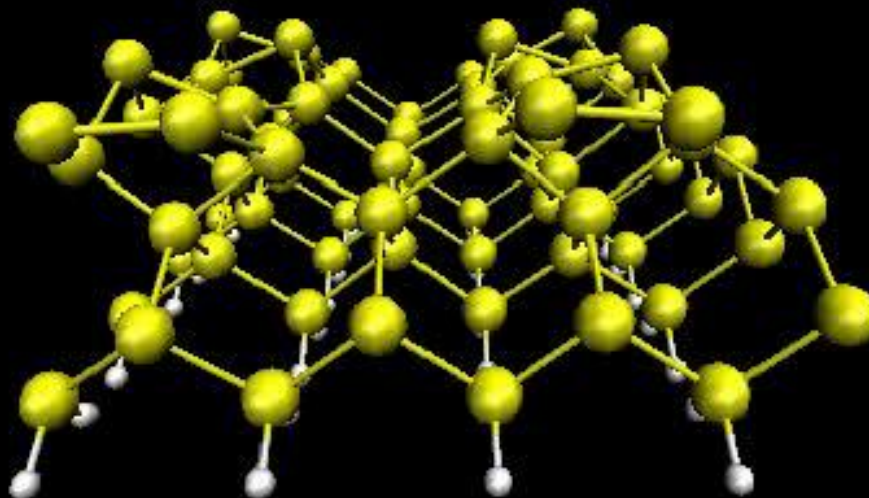
Non-concerted mechanism of Diels-Alder adduct from *ab initio* MD

P. Minary and MET *J. Am. Chem. Soc.* **126**, 13920 (2004)

P. Minary and MET *J. Am. Chem. Soc.* **127**, 1110 (2005)

R. L. Hayes and MET *J. Am. Chem. Soc.* **129**, 12172 (2007)

R. Iftimie, P. Minary, MET *Proc. Natl. Acad. Sci.* **102**, 6654 (2005)



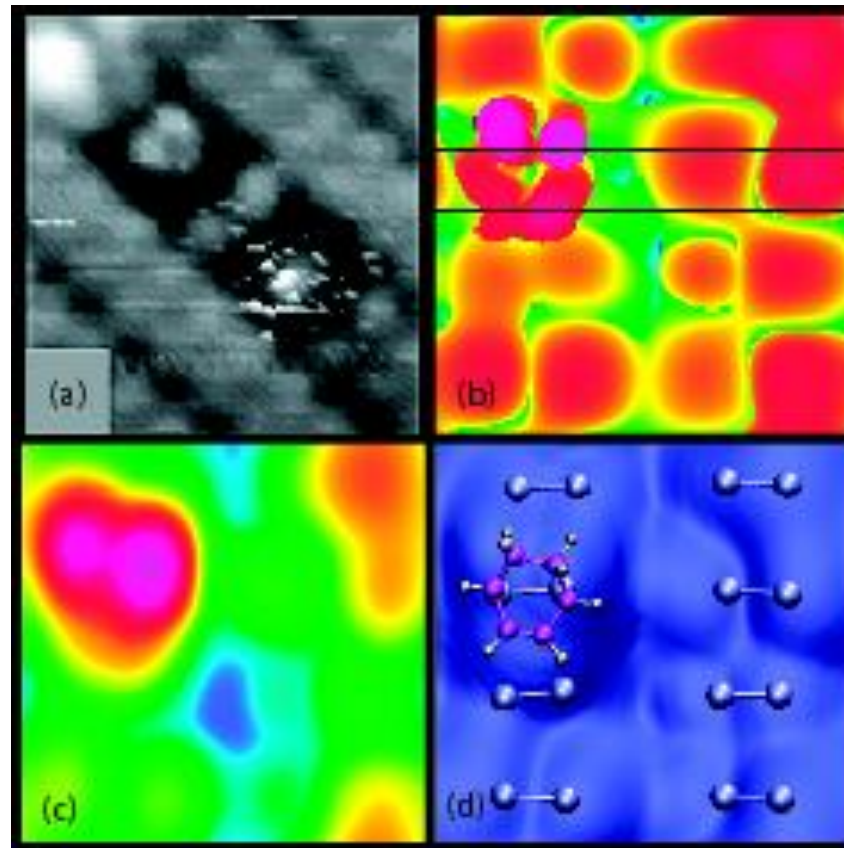
1,3-cyclohexadiene
on Si(100)-2x1

Calculation of Empty-State STM Images

R. L. Hayes and MET (to be submitted)

L. C. Teague and J. J. Boland
JPCB **107**, 3820 (2003)

Tersoff-Hamann



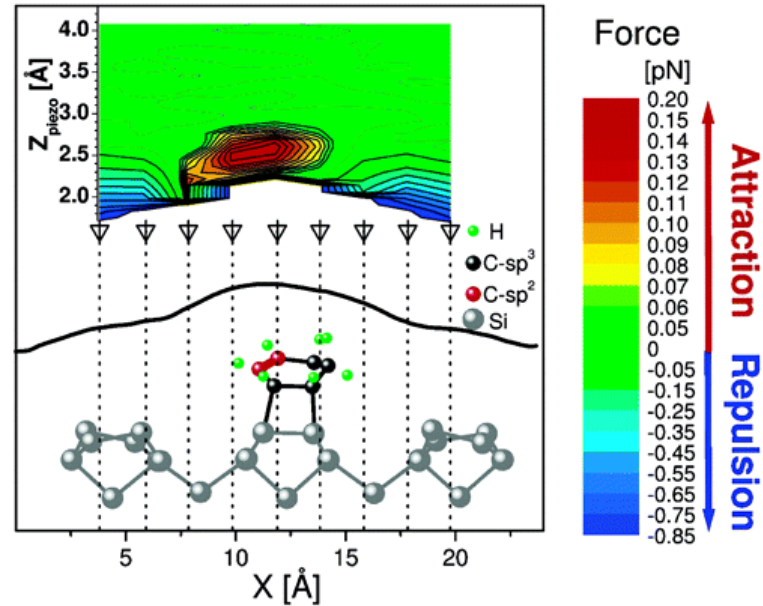
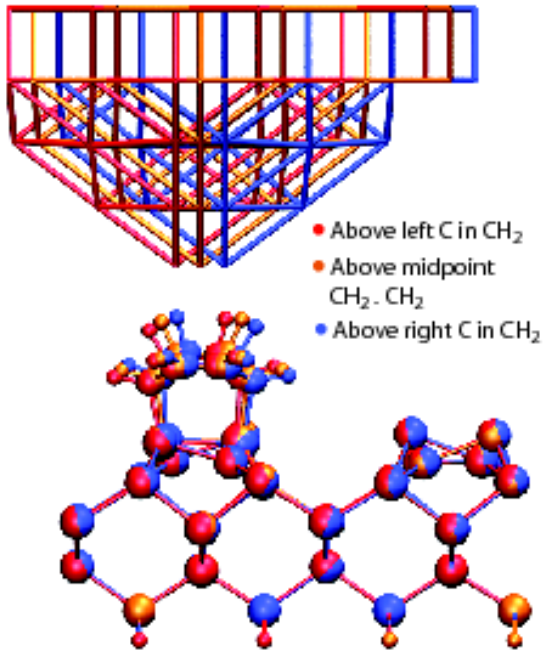
Bardeen

Tersoff-Hamann:
$$I \propto \sum_{\nu} \left| \psi_{\nu}(\mathbf{r}_{\text{tip}}) \right|^2 \delta(E_{\nu} - E_F)$$

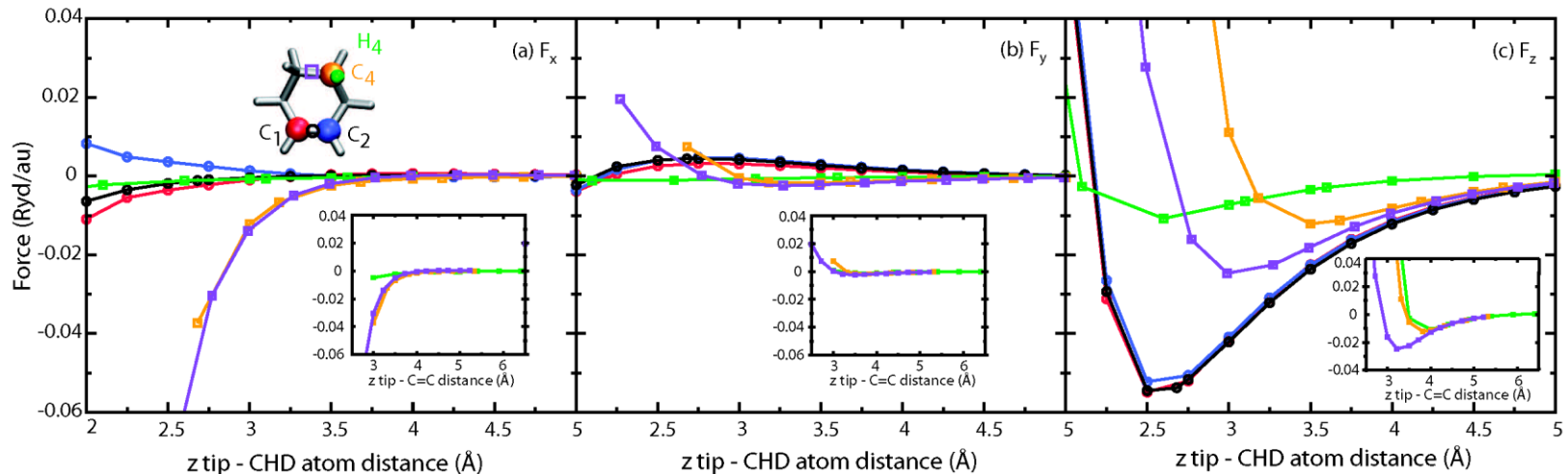
Bardeen:
$$I \propto \sum_{\mu, \nu} f(E_{\mu}) [1 - f(E_{\nu} + eV)] \left| M_{\mu\nu} \right|^2 \delta(E_{\mu} - E_{\nu})$$

Calculation of STM Images

Naydenov *et al. Nano Lett.* **6**, 1752 (2006)



Need to perform self-consistent conductance calculations (NEGF) on a full grid!



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Postdocs

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