

**October 22-24, 2003 CSCAMM Workshop:
Fundamental Physical Issues in Nonequilibrium Interface Dynam**

Drift-Induced Pattern Formation on Si(001) Vicinal Surfaces

- Effect of Alternating Anisotropy -

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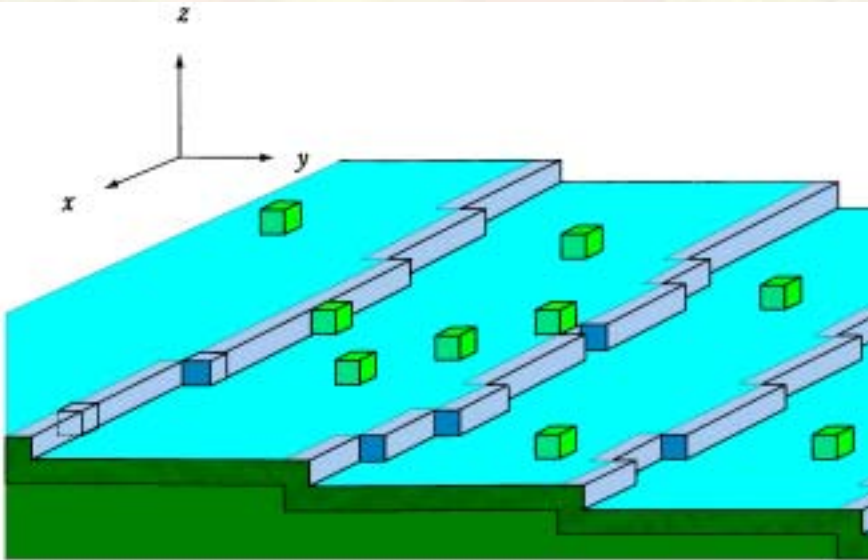
uwaha@phys.nagoya-u.ac.jp

Plan of My Talk

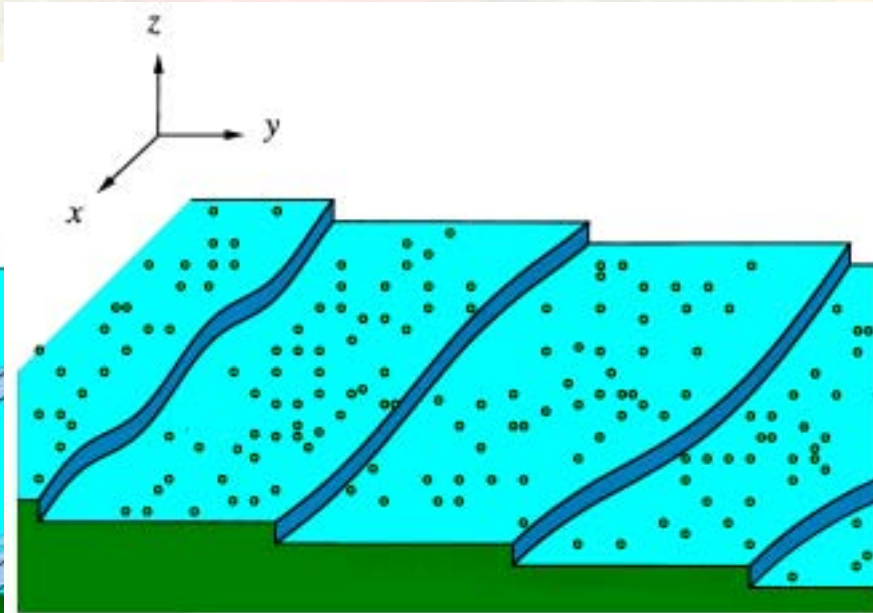
- Preliminaries: models, instabilities of a vicinal face---step bunching and step wandering on Si(111) vicinal faces.
- Characteristics of the Si(001) vicinal face and experimental observations of instabilities under direct current heating.
- Our theoretical study and Monte Carlo simulation: mechanism of step bunching and a new kind of wandering instability

Lattice model and continuum step model

Lattice model:
Monte Carlo simulation

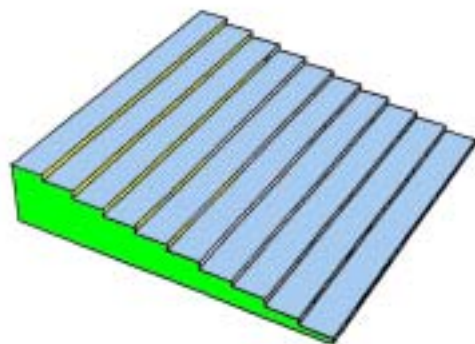


Continuum model:
Mathematical Analysis
Simulation in 1D

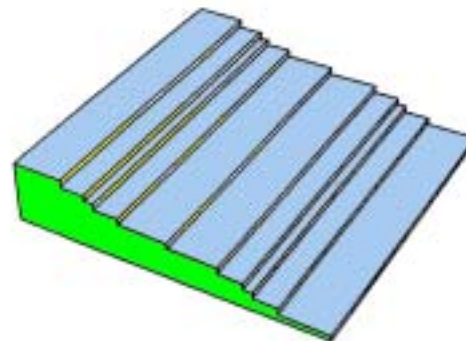


Morphological instabilities on a vicinal surface

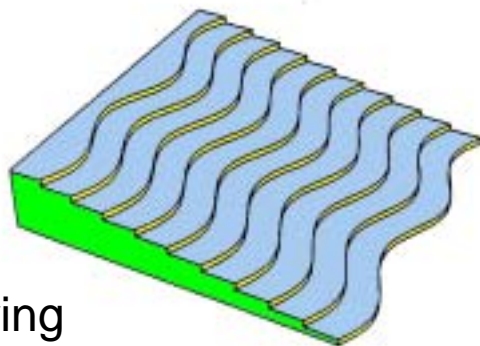
vicinal face



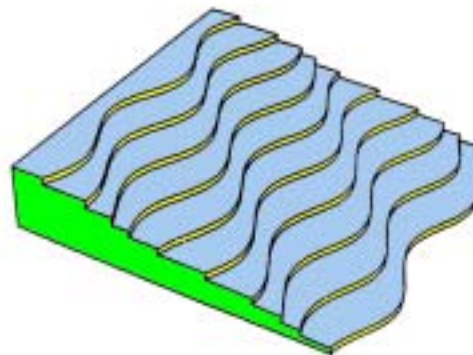
bunching



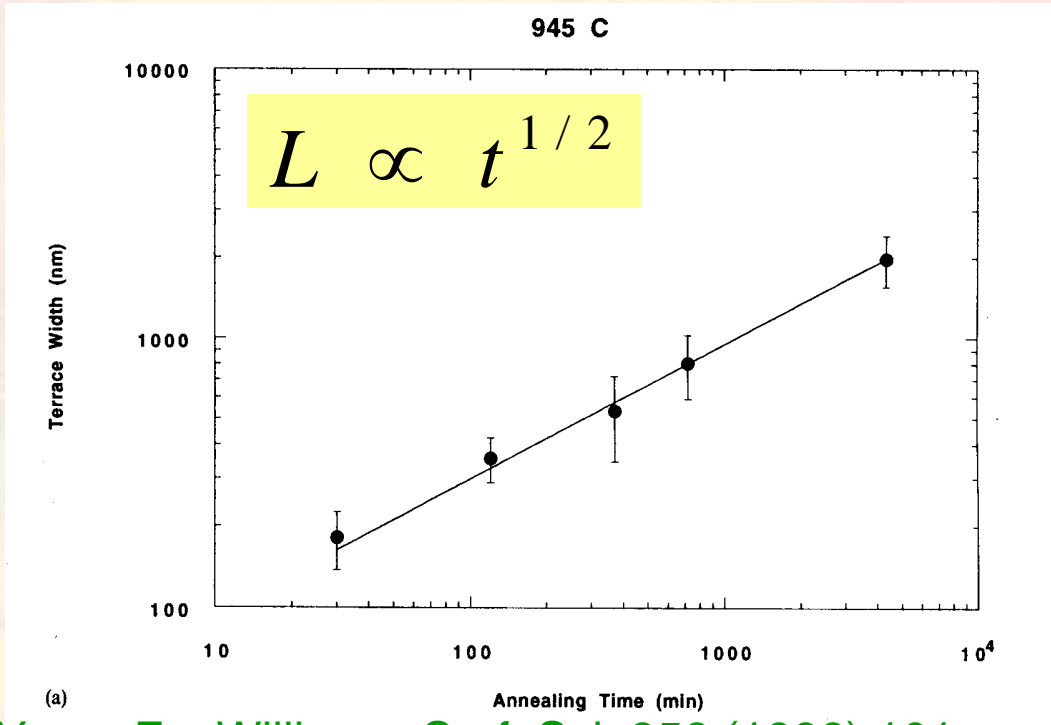
In-phase wandering



bending

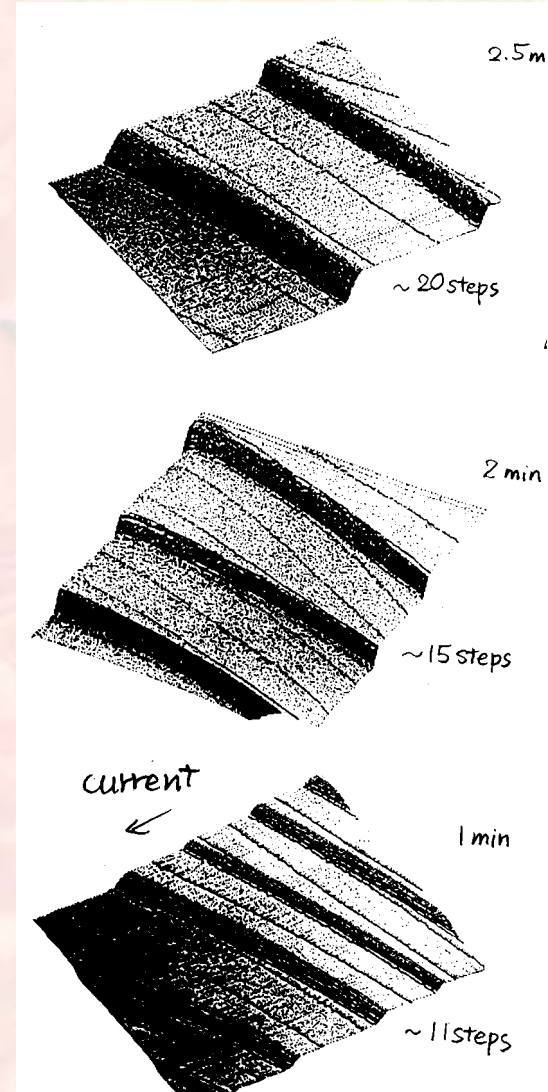


Step bunching on a Si(111) vicinal face under direct current heating



Yang, Fu, Williams: Surf. Sci. 356 (1996) 101

Fu, Liu, Johnson, Weeks, Williams:
Surf. Sci. 385 (1997) 259



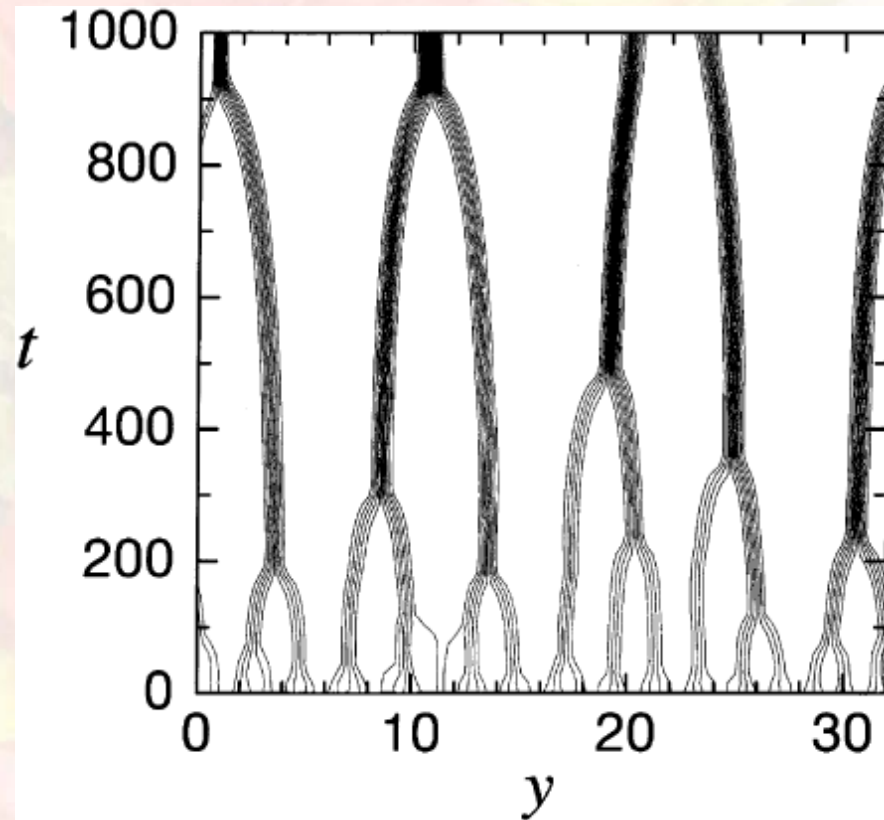
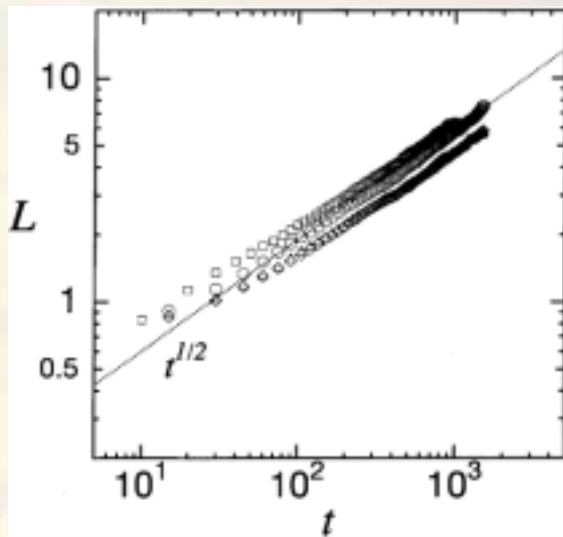
Theoretical study on growth of step bunches in Si(111) with drift

- C. Misbah, O. Pierre-Louis: Phys. Rev. E **53** (1996) R4319
- M. Sato, M. Uwaha: J. Phys. Soc. Jpn. **65** (1996) 1515
Step bunches (during evaporation) are described by Benney equation **Steady State of bunches.**
- H. Dobbs, J. Krug: J. Phys. I France **6** (1996) 413
MC simulation of SOS lattice model $t^{1/4}$ (1D), $t^{1/2}$ (2D)
- D.-J. Liu, J. D. Weeks: Phys. Rev. B **53** (1998) 14891
Continuum model + scaling hypothesis $t^{1/2}$ growth(1D)

Step bunching on a Si(111) vicinal face under direct current heating—1D model

- Instability with step-down drift
Formation of step pairs
- Equidistant pairs are also unstable Hierarchical pairing
- Scaling laws in bunching

$$L = t^{1/2}$$

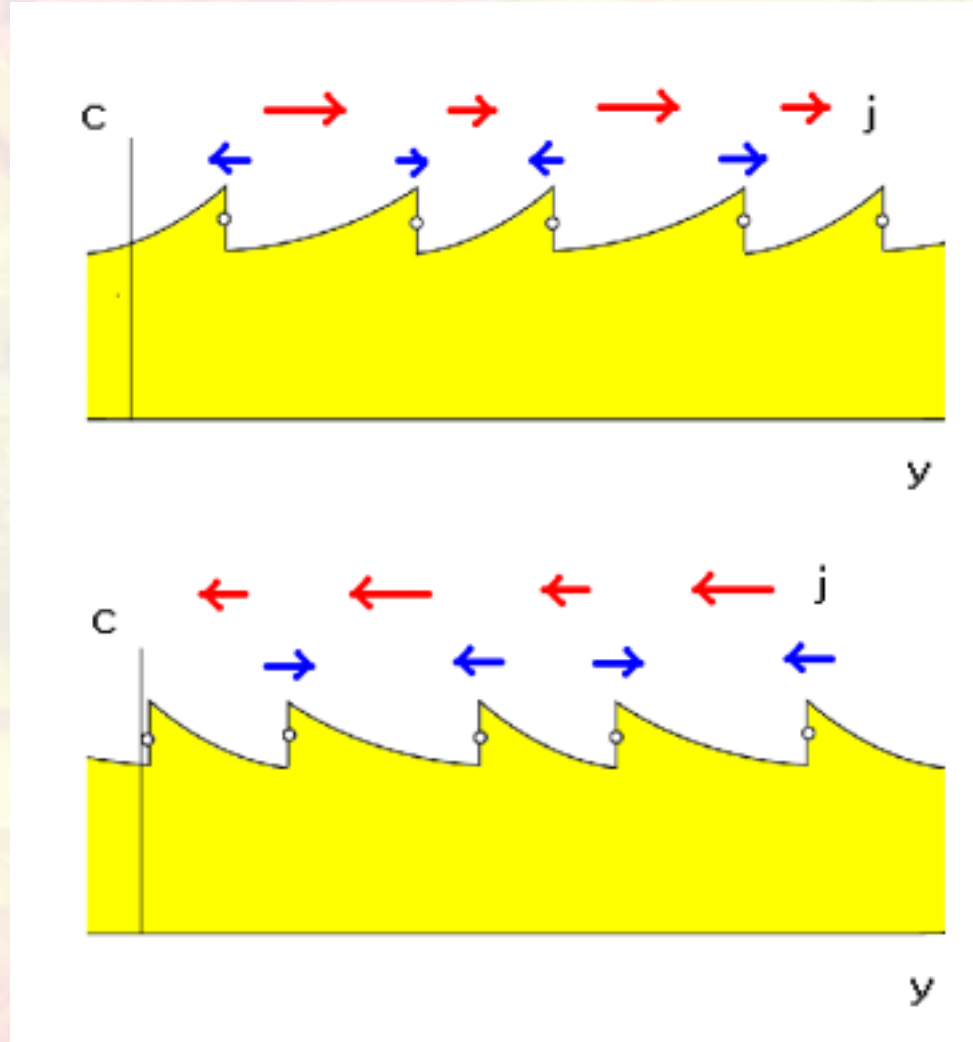


Mechanism of step pairing in a Si(111) vicinal

Current on a terrace

$$j \approx -D_s \frac{c_- - c_+}{L} + \frac{c_- + c_+}{2} v_d$$

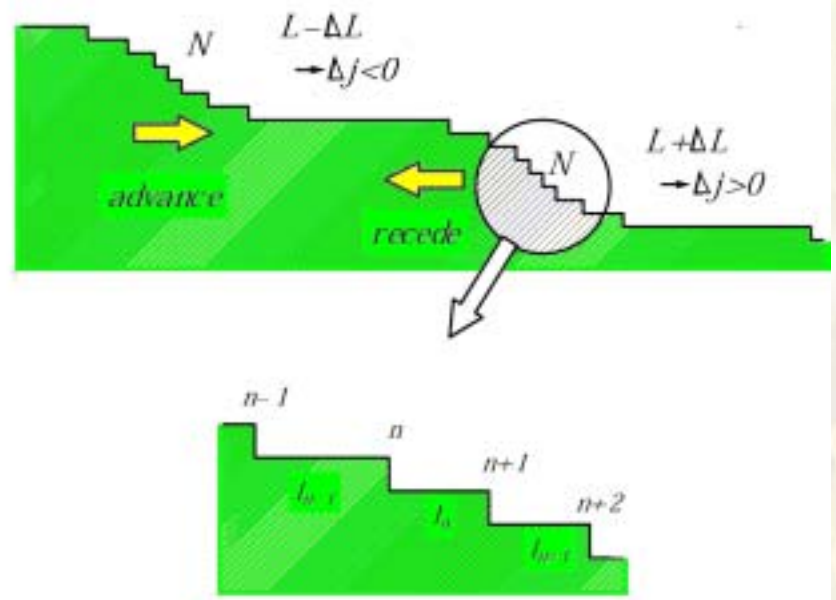
Density profile of adatoms



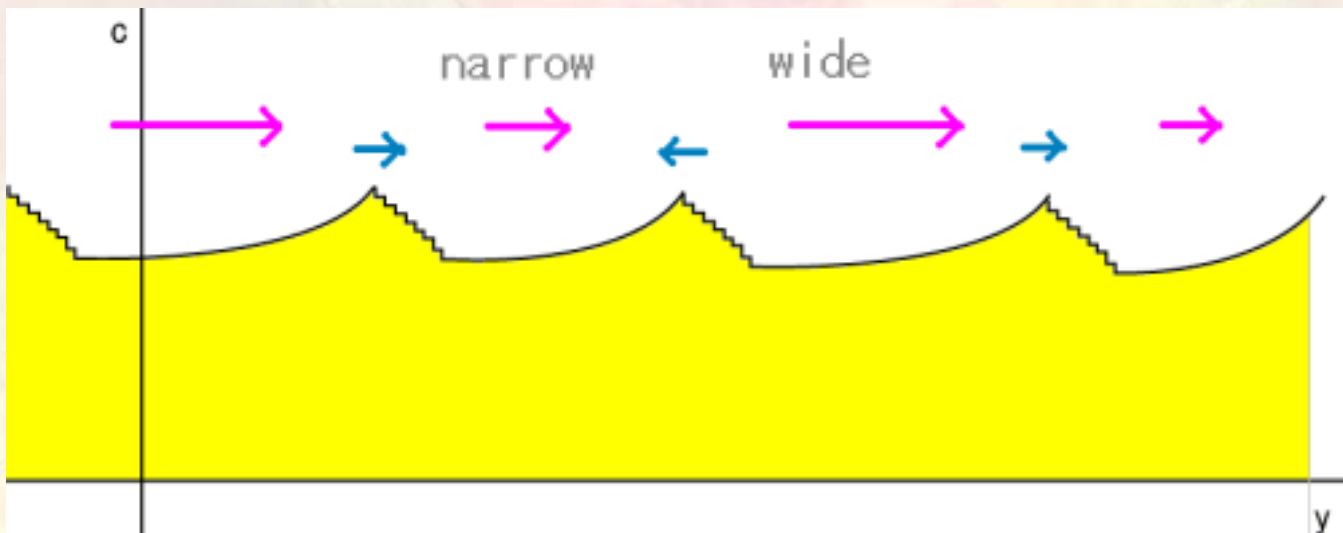
Mechanism of bunch coalescence on a Si(111) vicinal

Current on a terrace

$$j \approx -D_s \frac{c_1 - c_N}{L} + \frac{c_1 + c_N}{2} v_d$$



Density profile of adatoms



Origin of the growth law of bunches

Current on a terrace

$$j \approx -D_s \frac{c_1 - c_N}{L} + \frac{c_1 + c_N}{2} v_d$$

Change of the terrace width

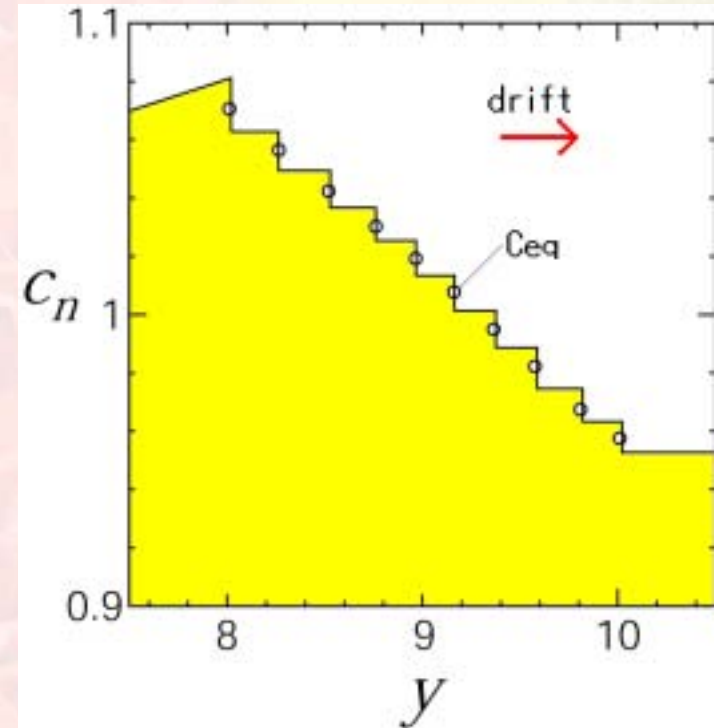
$$\frac{d\Delta L}{dt} \propto \frac{\Delta j}{N} \propto \frac{1}{N} \frac{c_1 - c_N}{L^2} \Delta L$$

Collision time

$$\tau \propto \frac{L^2}{v_d} \propto N^2$$

One can show $c_1 - c_N \propto N$

And $l_b \propto N^{-2/(\nu+1)}$

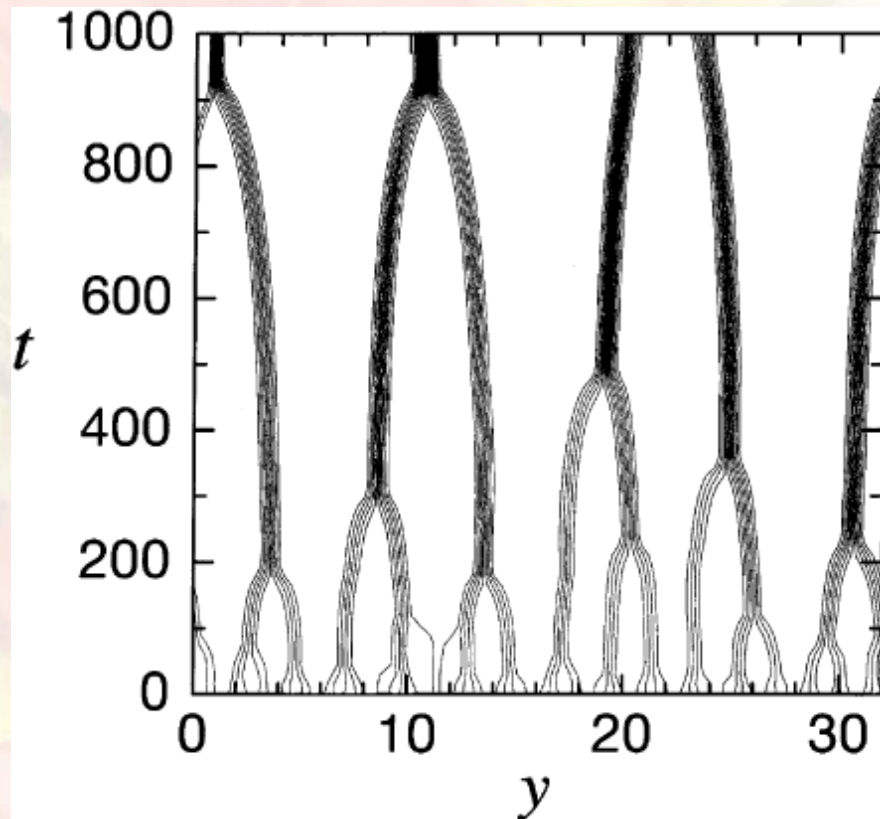
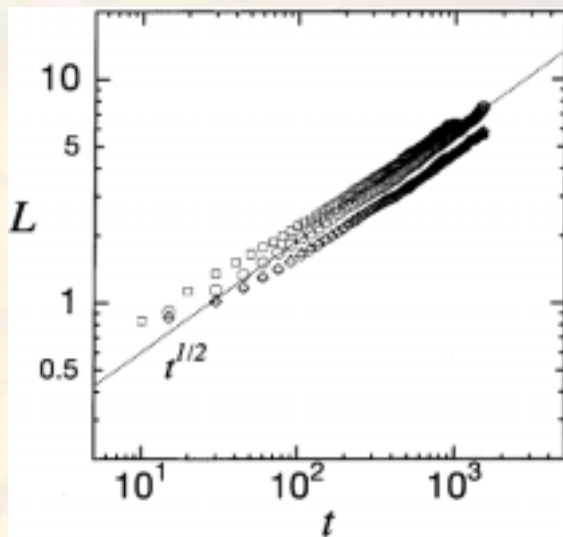


Step repulsion potential I

Step bunching on a Si(111) vicinal face under direct current heating—1D model

- Instability with step-down drift
Formation of step pairs
- Equidistant pairs are also unstable Hierarchical pairing
- Scaling laws in bunching

$$L = t^{1/2}$$



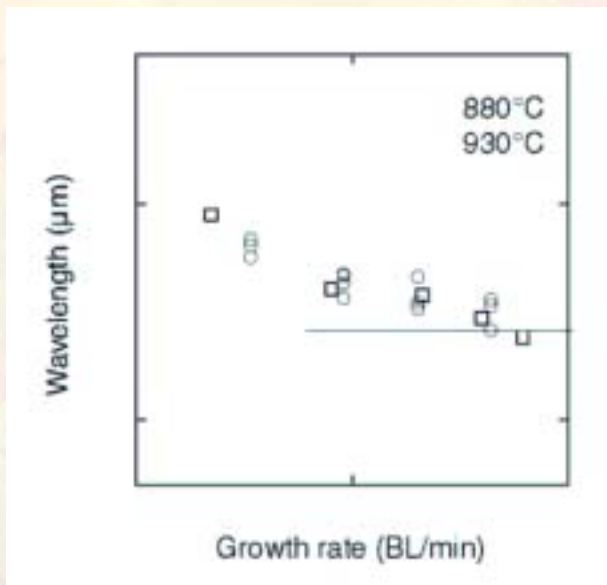
Causes of step wandering instability

Asymmetry of the diffusion field in front and in back of the step

Mechanisms

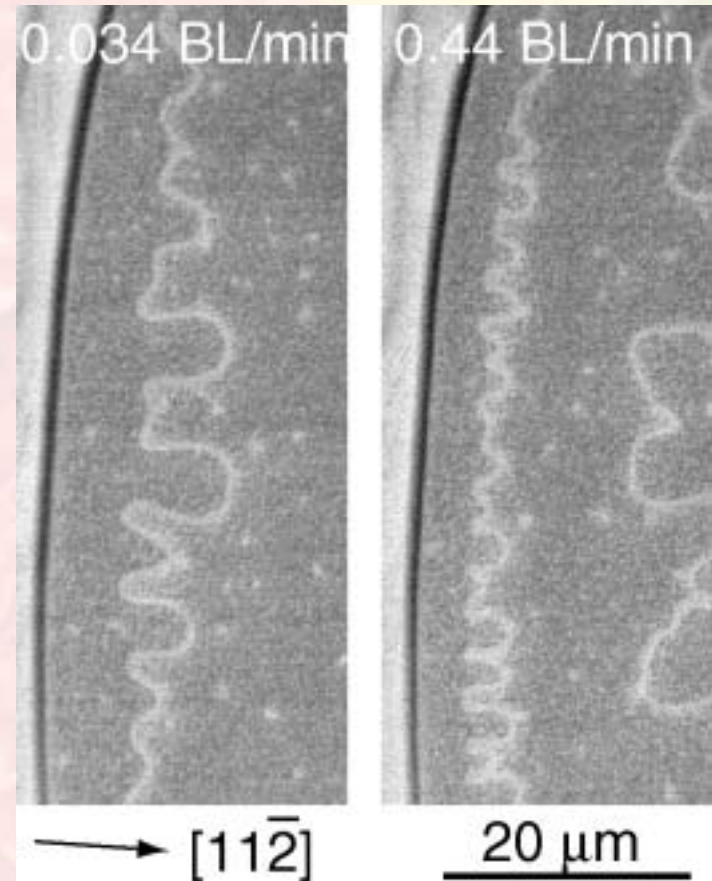
- Ehrlich-Schwoebel (ES) effect (Bales, Zangwill 1990)
- Kink Schwobel effect (Pierre-Louis, D'Orsogna, Einstein 1999)
- Drift flow of adatoms by an external field (Sato, Uwaha 1996)
- Difference of the terrace width in front and in back of the step

Step wandering due the difference of the terrace width



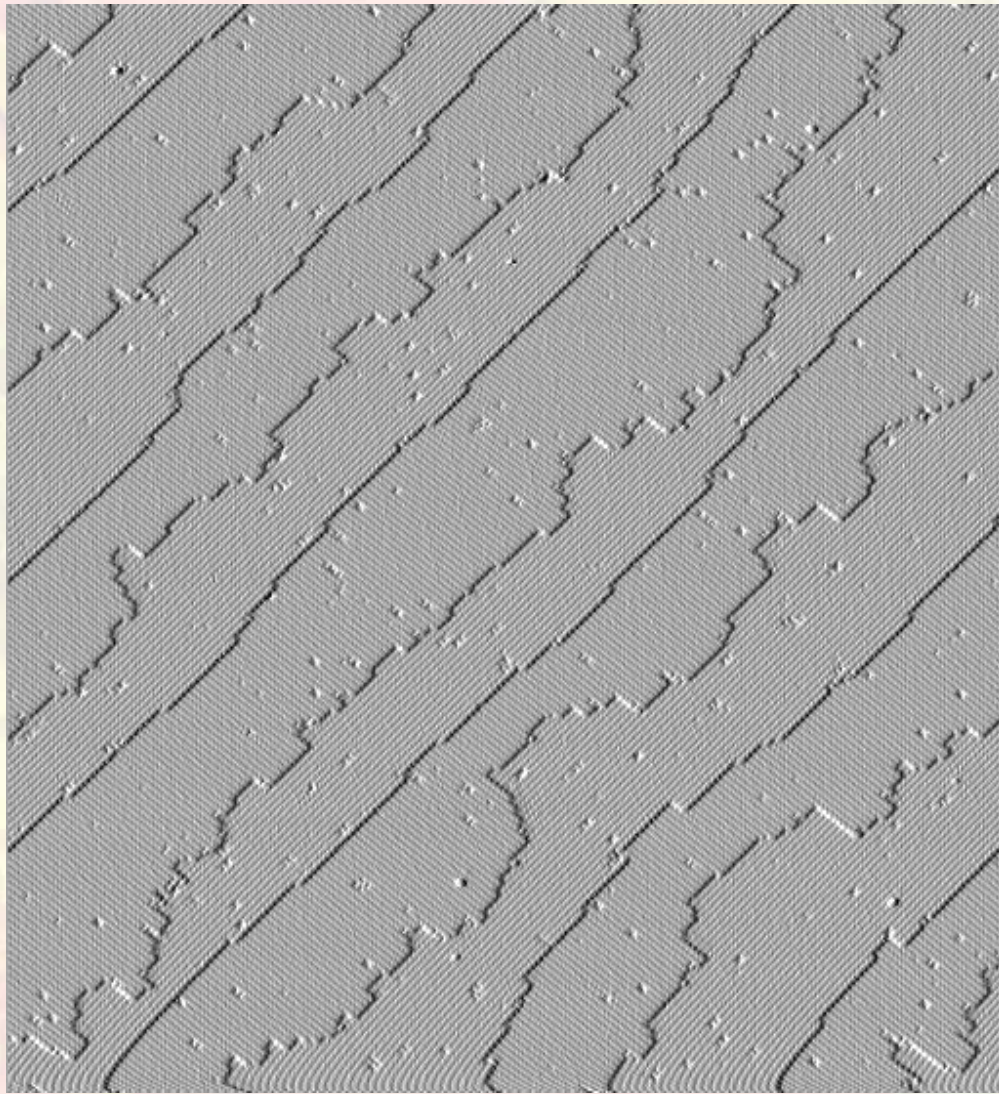
- Wandering wave number in the one-sided model

$$k_{\max} = \sqrt{\frac{k_B T V_0}{3 \Omega^2 \tilde{\beta} c_{eq}^0 D_s}}$$



Two kinds of steps and terraces on a Si(001) vicinal face

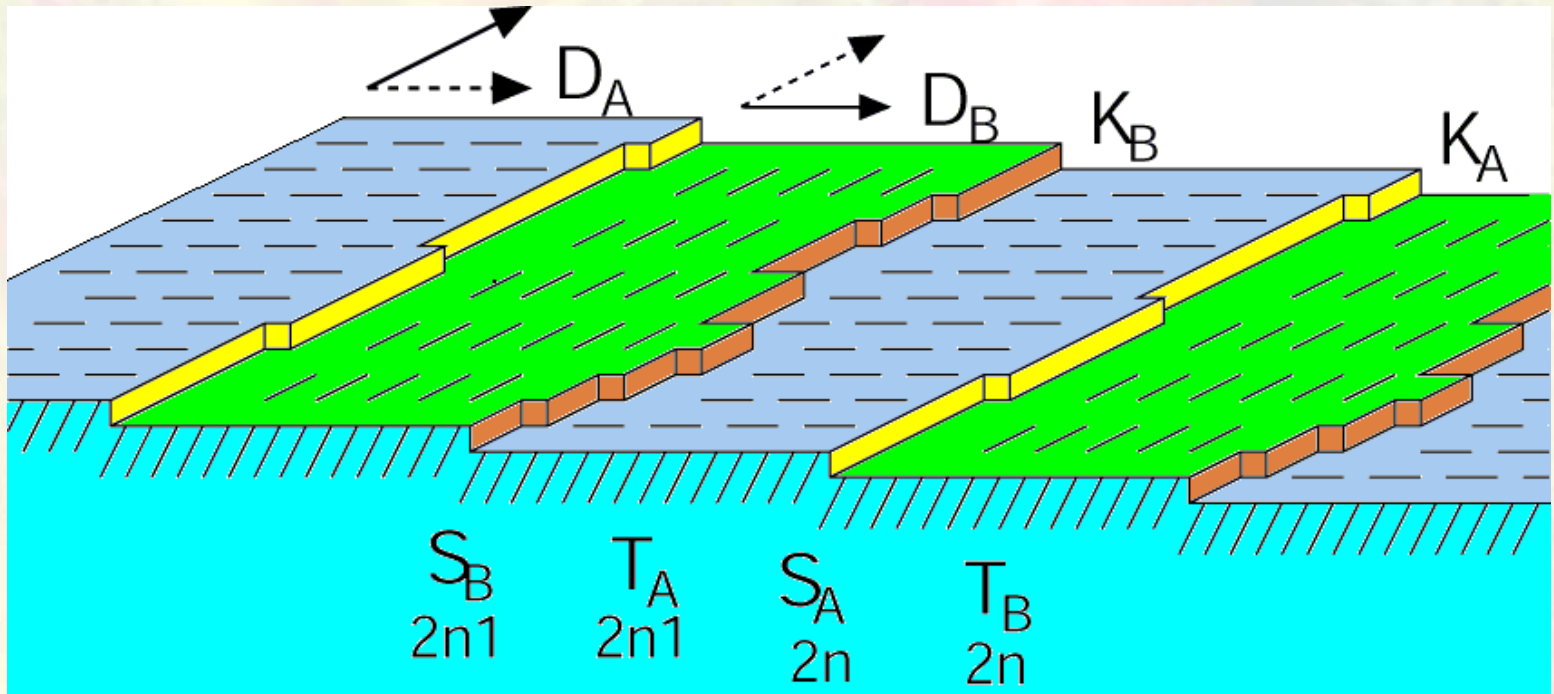
- STM image of a vicinal tilted towards [110]
- A-step with few kinks
- B-step with many kinks
- Stripes are dimer rows



(STM image taken by
Swartzentruber et al.)

Si(001) vicinal face tilted to [110]

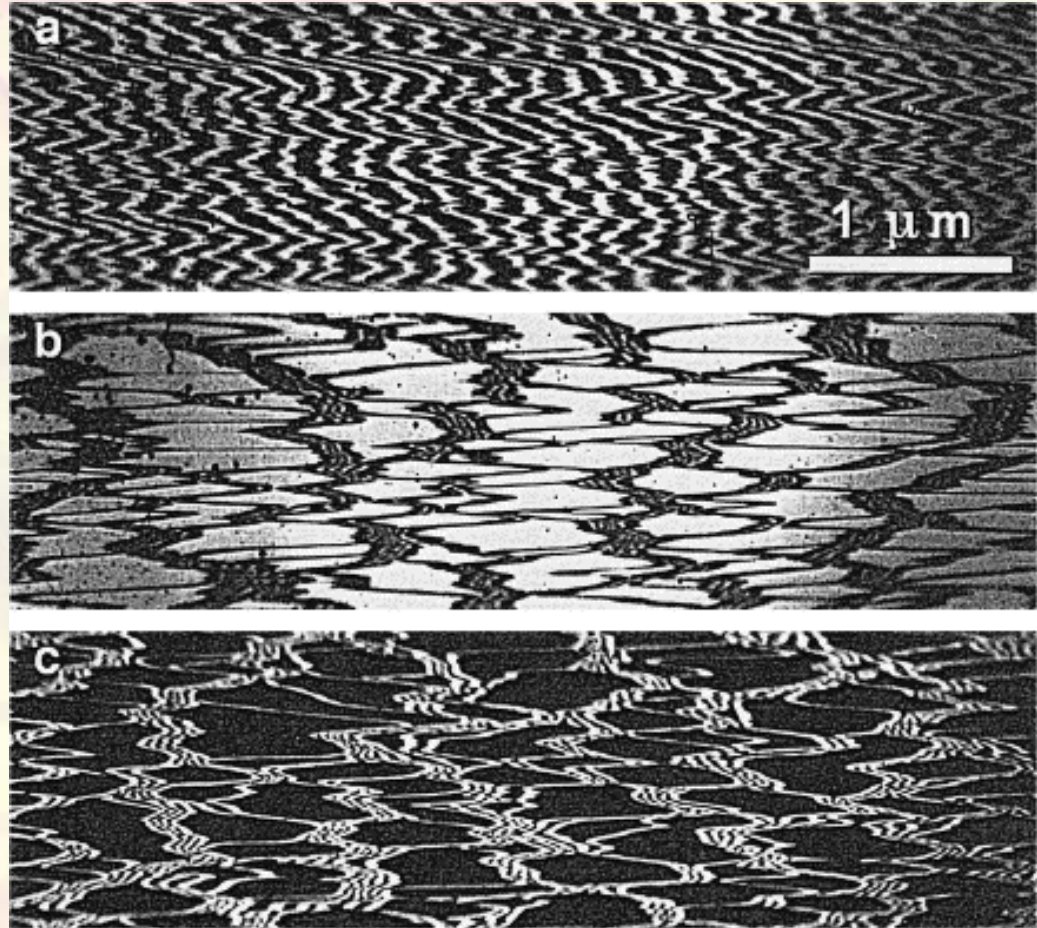
- The orientation of dimer rows alternates on neighboring terraces .



Si(001) step pairing with drift

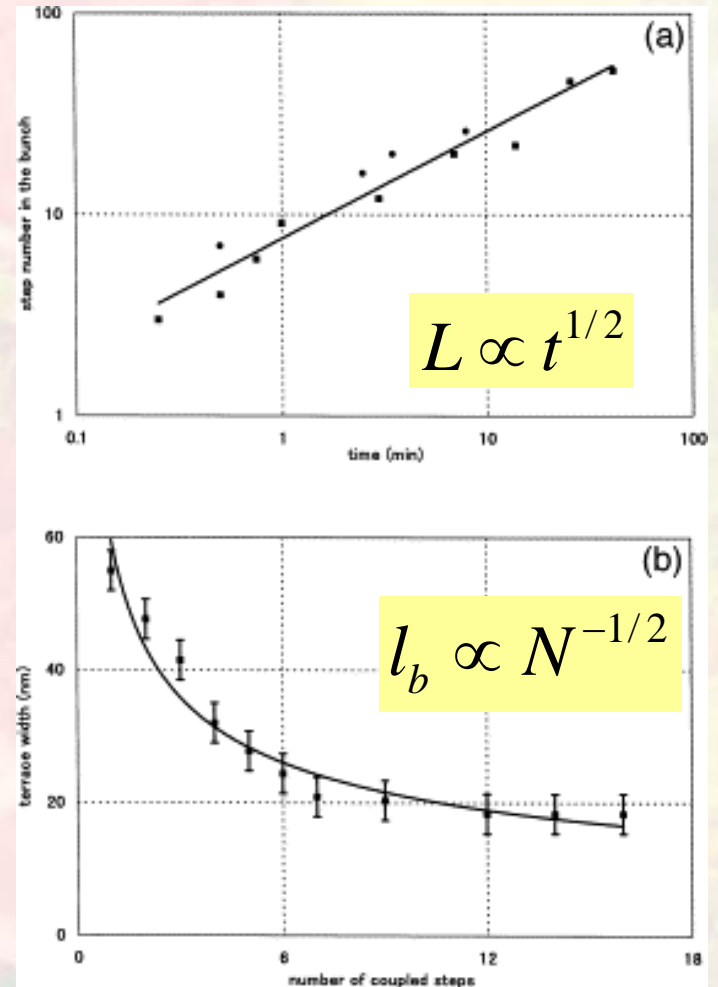
REM images of the stepped Si(001) surface.

- (a) After AC heating.
- The same surface after DC heating at 1150 ° C for 2 min:
(b) in the step-up,
(c) in the step-down directions



Si(001) step bunching with drift

- (a) Change of the average distance between bunches. The power-law with an exponent 0.5.
- (b) Average distance between the steps in a bunch vs. the number of steps in the bunch. The exponent is -0.5.

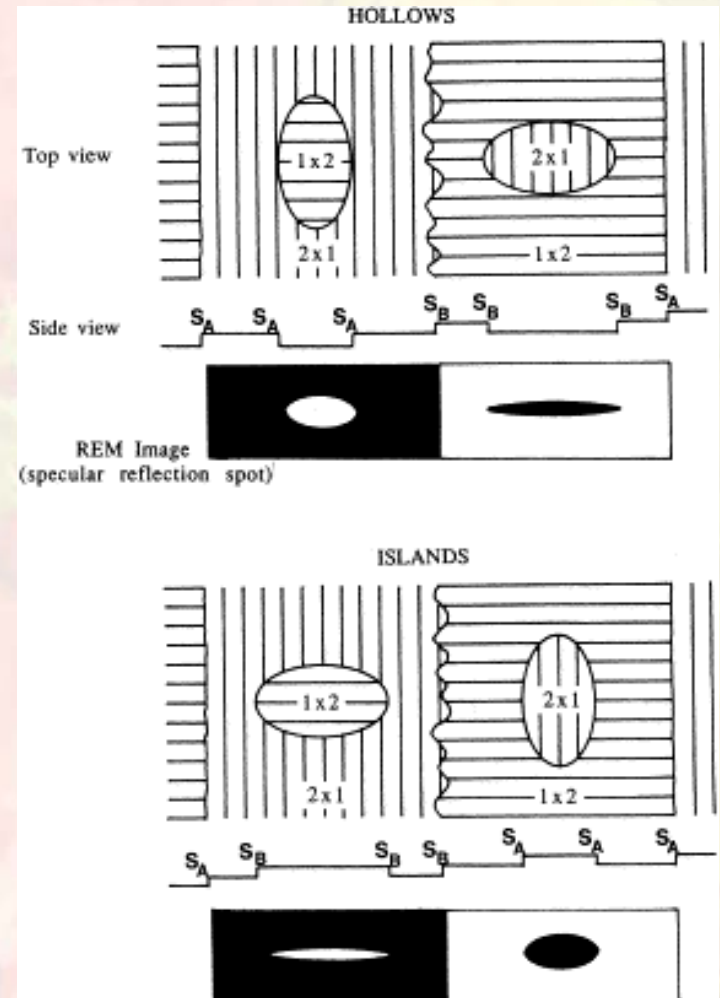


Si(001) island motion with drift

A stepped Si(001) surface with hollows and islands.

- The analysis of the island motion is consistent with the **positive** effective charge of adatoms.

Métois, Heyraud, Pimpinelli:
Surf. Sci. **420** (1999) 250



Theoretical study on bunching in Si(001)

- S. Stoyanov: Jpn. J. Appl. Phys. **29** (1990) L659
1D model Formation of step pairs due to the difference in adatom mobility (diffusion coefficient)
- A. Natori, H. Fujimura, H. Yasunaga: Jpn. J. Appl. Phys. **31** (1992) 1164
Numerical simulation of 1D model (repulsion, impermeable steps)
Formation of large bunches with step-up drift only.
- M. Sato, M. Uwaha, and Y. Saito: J. Cryst. Growth **237-239** (2002) 43
Stability analysis of 1D model . Monte Carlo simulation of 2D lattice model → bunching in both drift direction

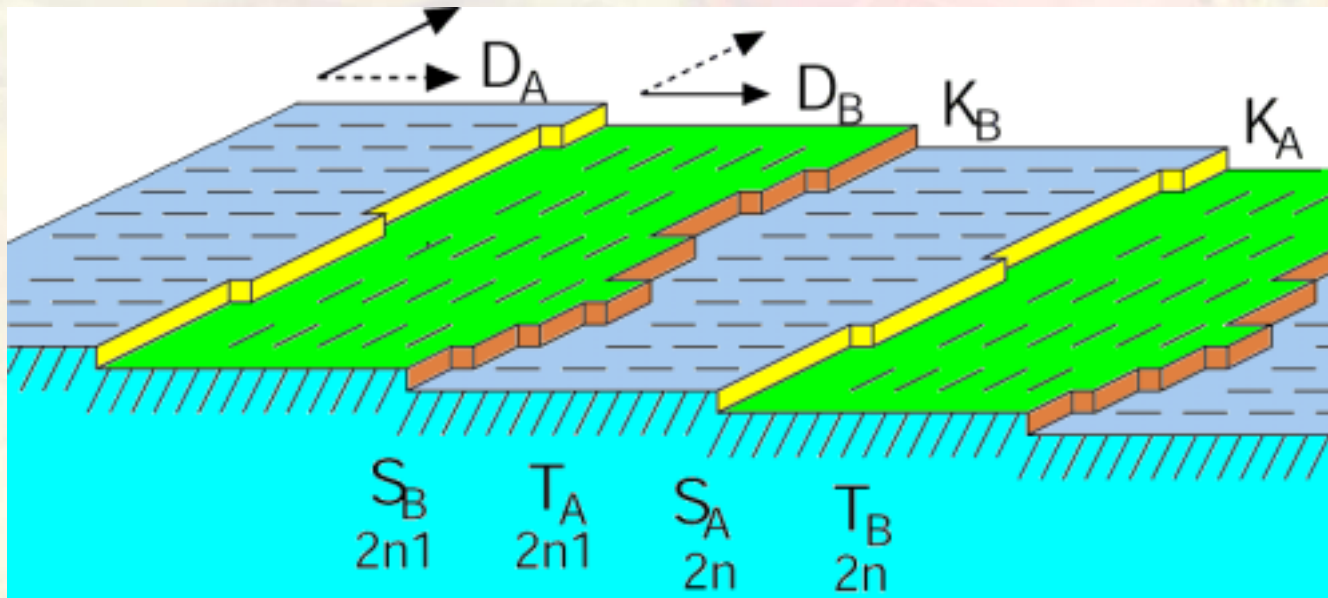
Vicinal face of Si(001)

- Alternating dimer row direction on consecutive terraces
Diffusion coefficients and drift velocity alternate .

$$D_B > D_A$$

$$v_{A,B} = \frac{D_{A,B}}{k_B T} Z_{eff} e E$$

Assume fast step kinetics



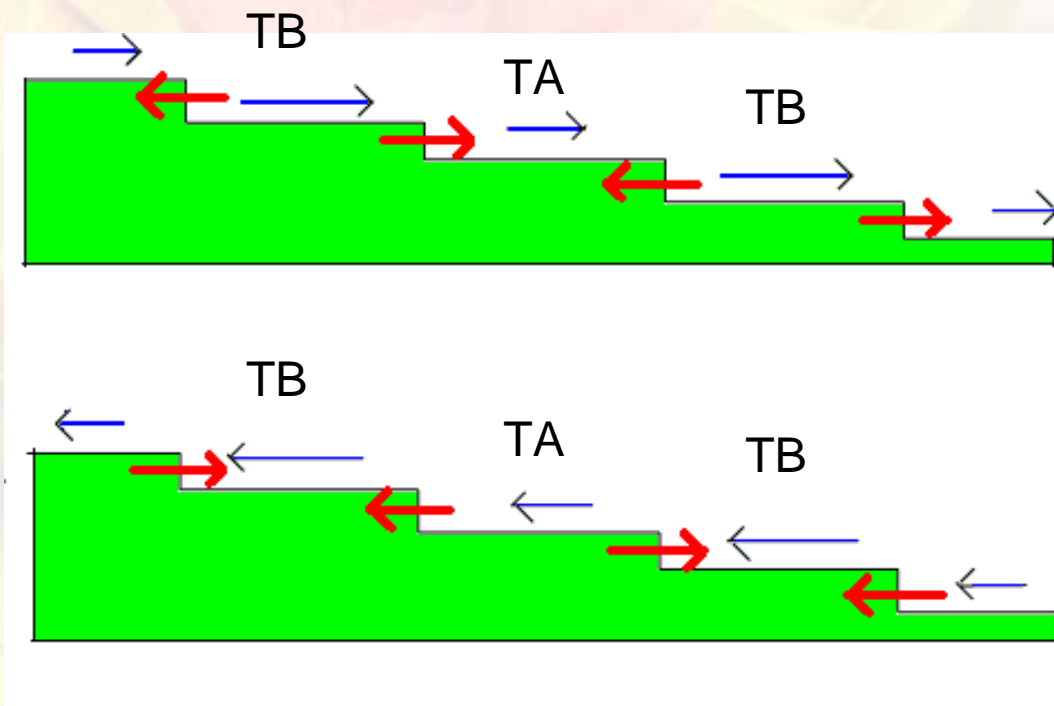
Formation of step pairs by the drift of adatoms on a Si(001) vicinal face (Stoyanov 1990)

- Step-down drift

Fast TB terraces expand

- Step-up drift

Slow TA terraces expand



Pairing of steps Formation of large bunches?

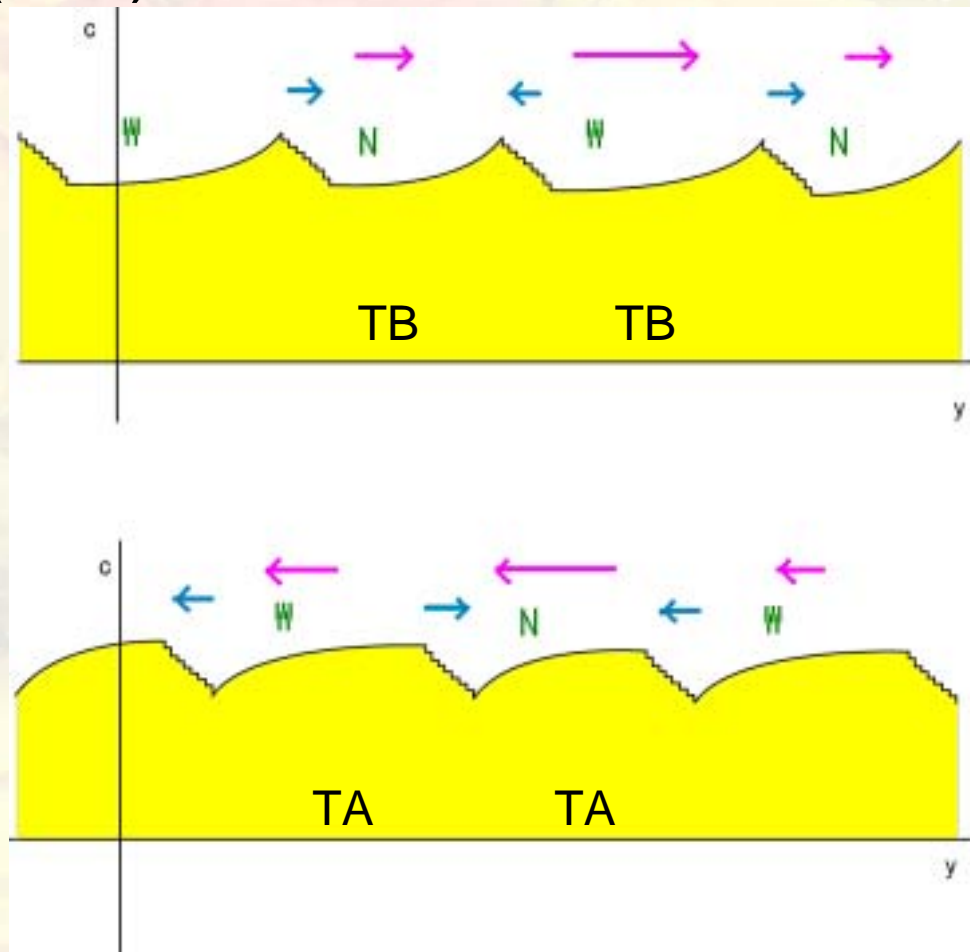
Formation of large bunches by drift of adatoms on a Si(001) vicinal face

- Step-down drift

Fast TB terraces expand

- Step-up drift

Slow TA terraces expand

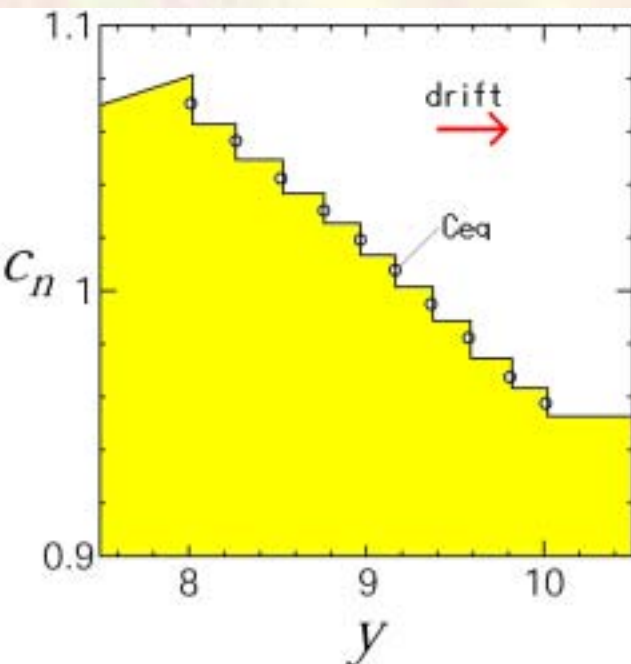


Pairing of steps

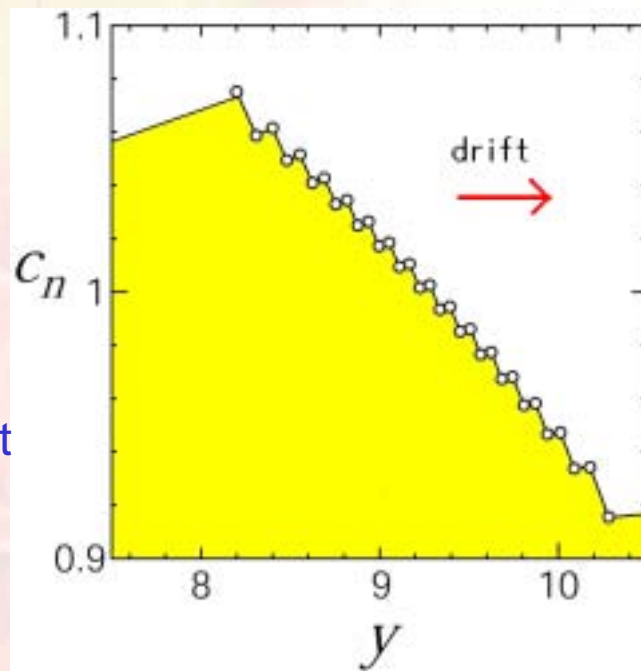
Formation of large bunches

Density of adatoms in a bunch

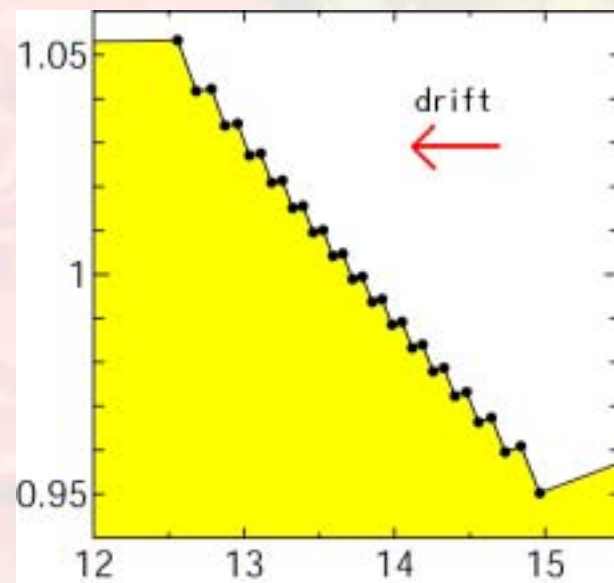
Si(111) step-down drift



Si(001)
step-down drift

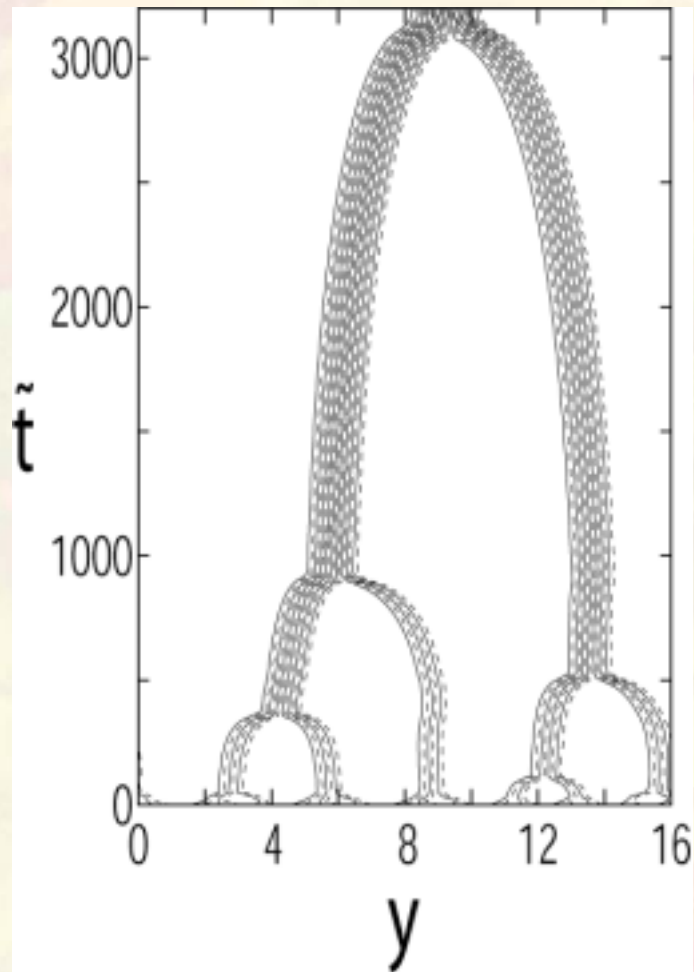


Si(001)
step-up drift

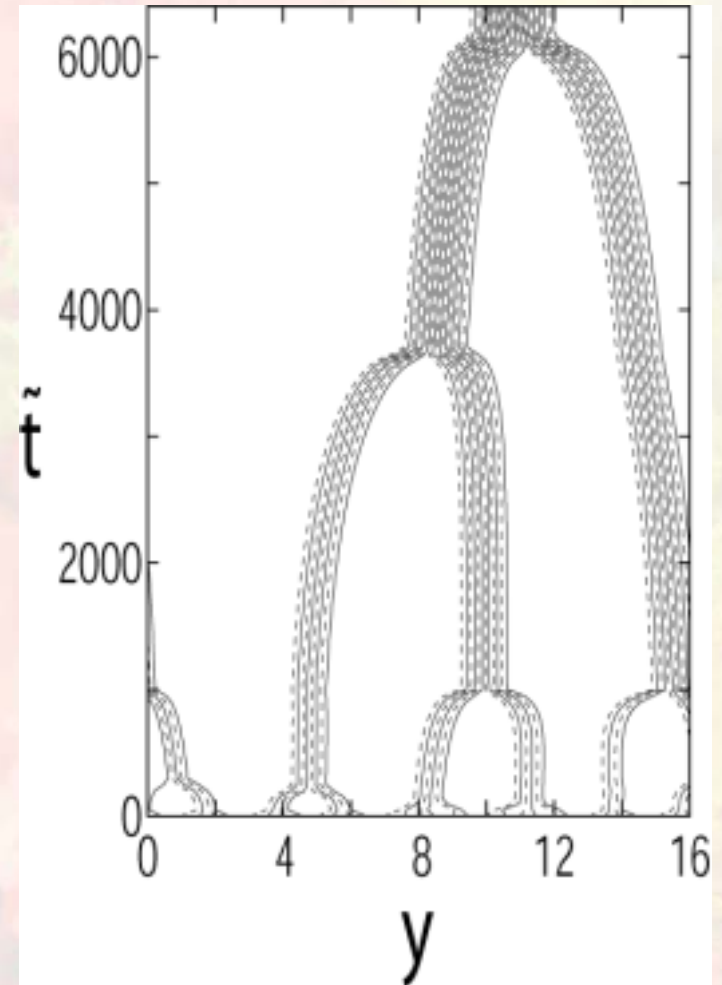


Drift direction and bunching Instability

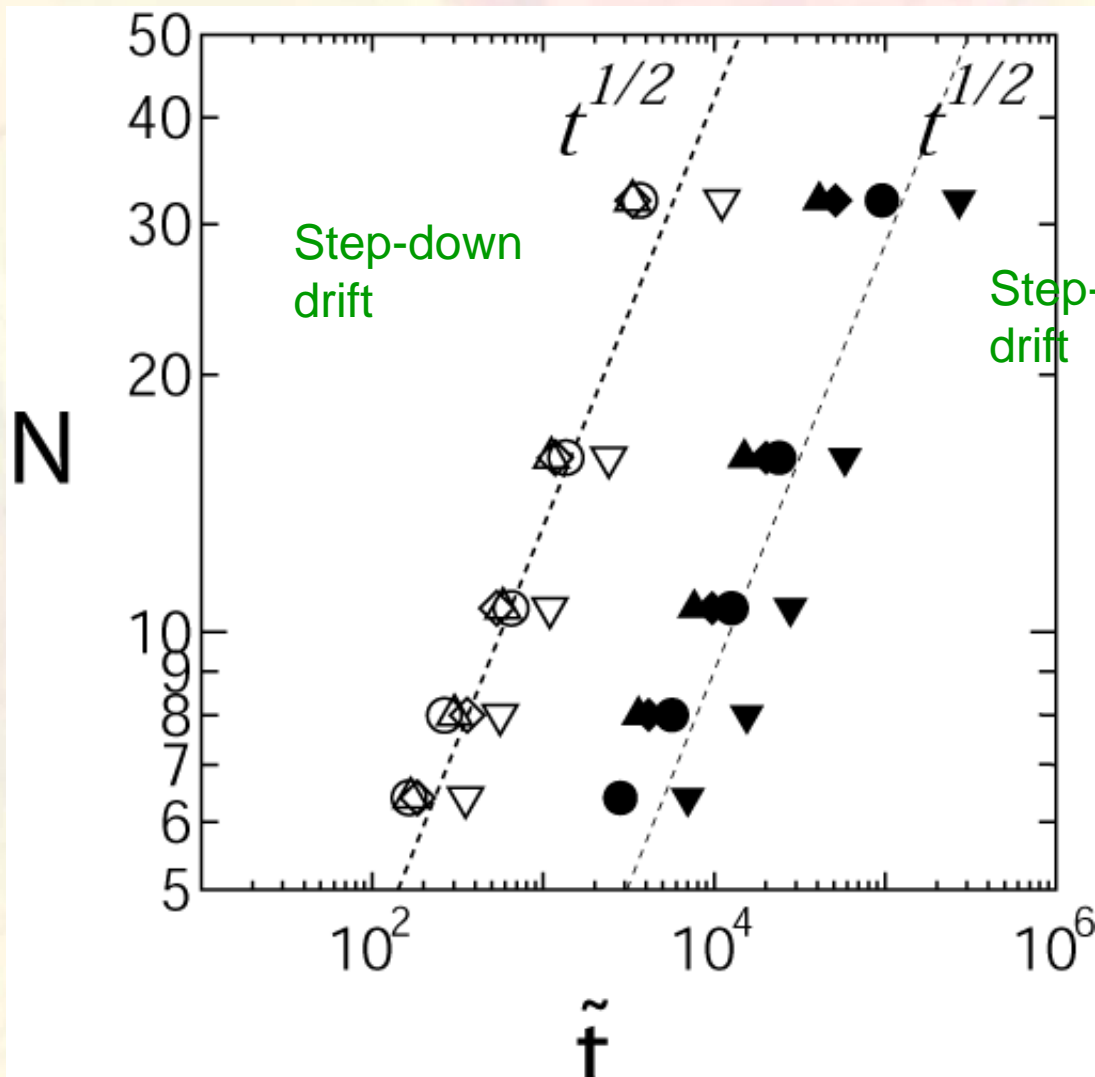
Step-down



Step-up



Growth rate of step bunches

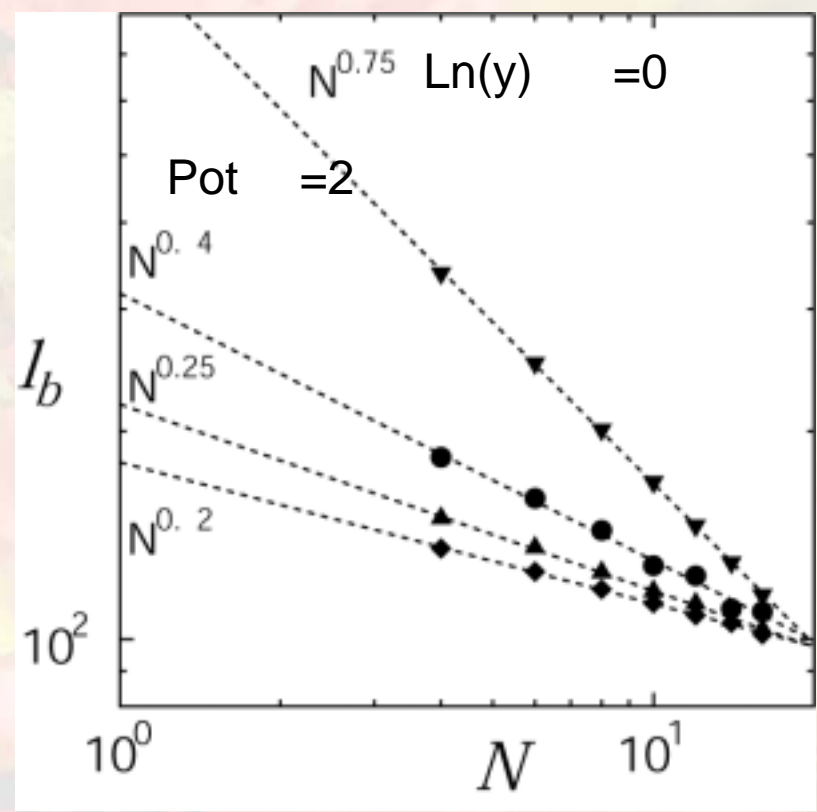
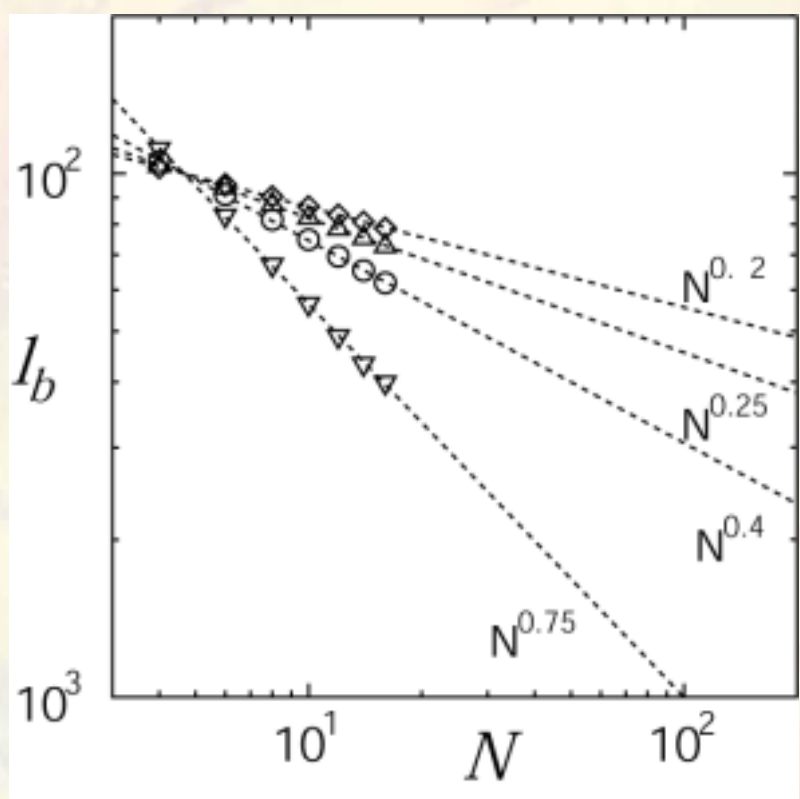


Step distance l_b in a bunch of size N with step repulsion potential $V(r) \sim |r|^{-\alpha}$

Step-down

Exponent = $3/2(\alpha + 2)$

Step-up



Anisotropic diffusion and drift in the Monte Carlo simulation

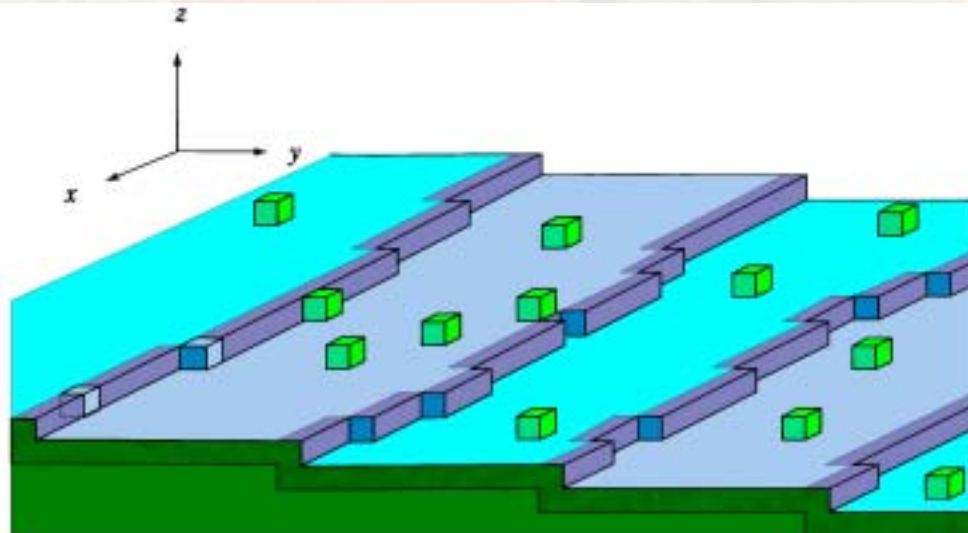
- Choose an atom

Terrace: A or B

Reduction factor for diffusion if the motion is in the slow direction

Drift is taken account as a biased diffusion.

$$\text{Pr}_{move} = \frac{1}{4} p \left(1 \pm \frac{Z_{eff} e E a}{2k_B T} \right)$$



Drift direction and bunching instability

Step-down

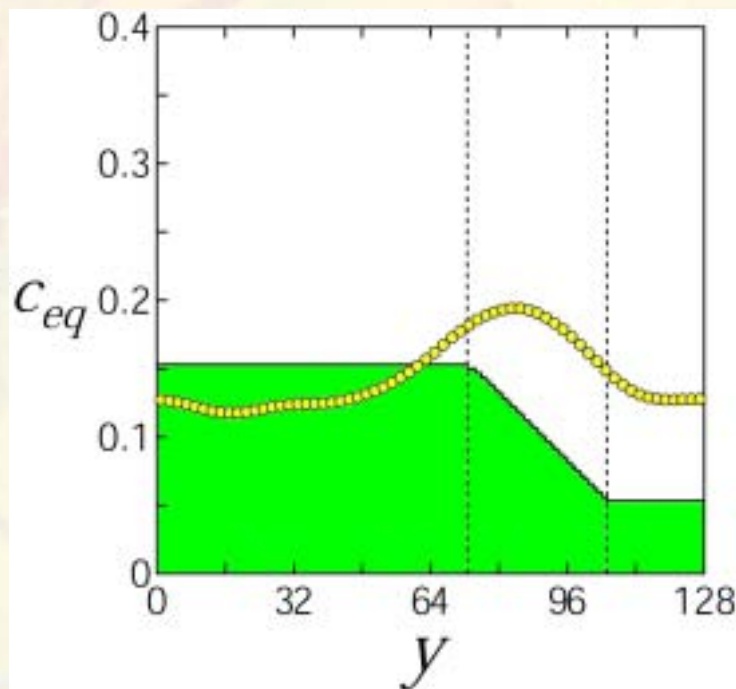
Step-up

Heavy figures deleted. See

M. Sato, M. Uwaha, and Y. Saito: *J. Cryst. Growth* **237-239** (2002) 43

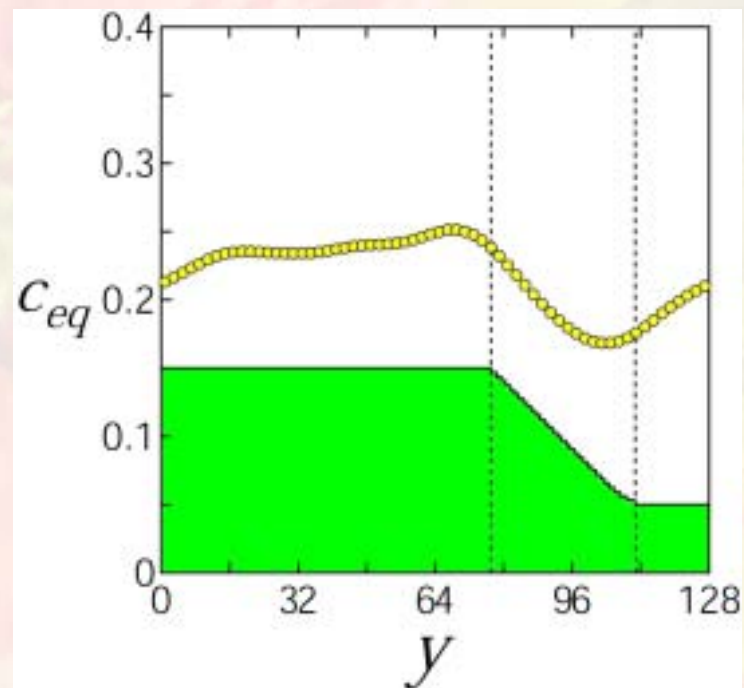
Average profile of the surface and adatom density

Step-down



Drift

Step-up



Drift

Drift direction and growth rate of bunches

Step-down

$$N_{bunch} \propto t^{1/2}$$

Step-up

Heavy figures deleted . See

M. Sato, M. Uwaha, and Y. Saito: J. Cryst. Growth **237-239** (2002) 43

Drift direction and bunching instability

Step-down

Step-up

Heavy figures deleted . See

M. Sato, M. Uwaha, and Y. Saito: J. Cryst. Growth **237-239** (2002) 43

Drift direction and instability in MC simulation with strong step repulsion

Step down

Step up

Heavy figures deleted. See

M. Sato, M. Uwaha, Y. Saito and Y. Hirose, Phys. Rev. B 67, 125408 (2003)

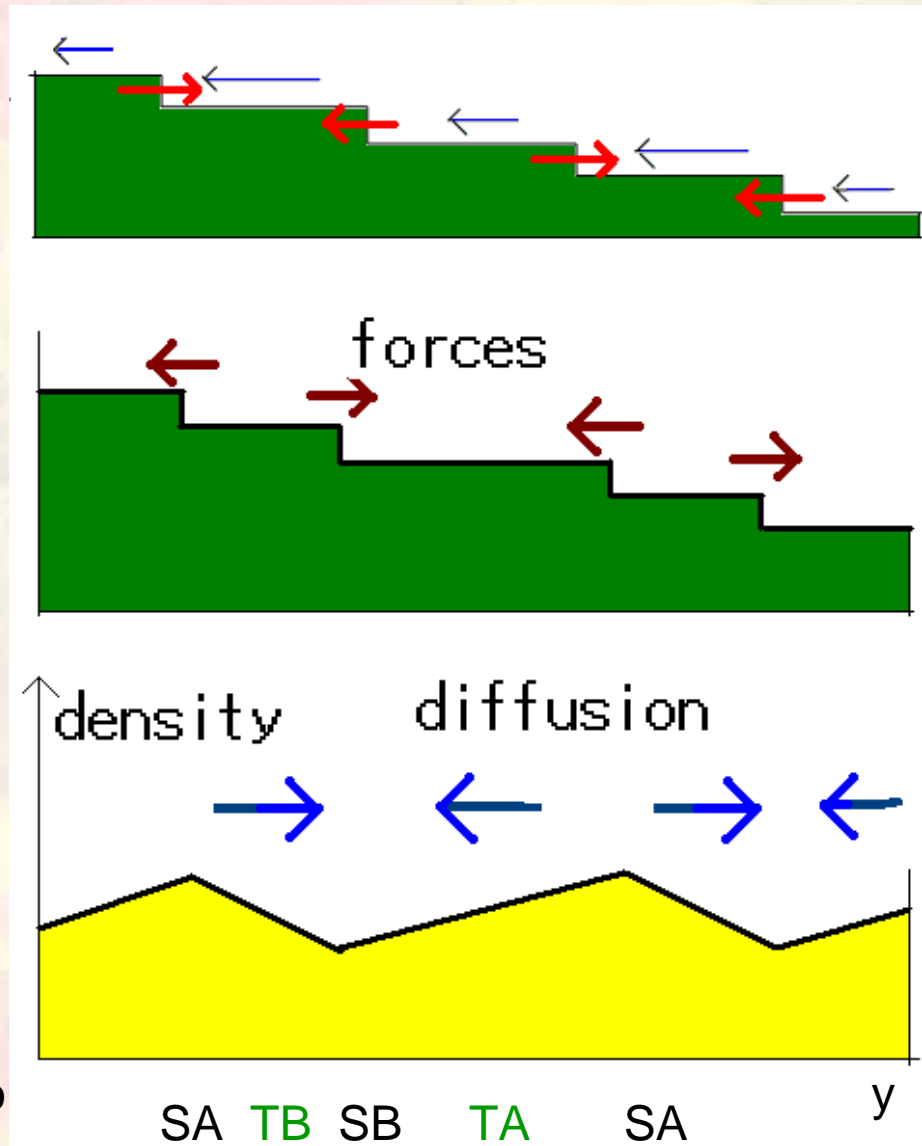
Steady state of step pairs with step-up drift

- Step-up drift

Slow TB terraces expand

Repulsion produces alternating density gradient

Diffusion current perpendicular to the drift is **large on slow TB terraces**



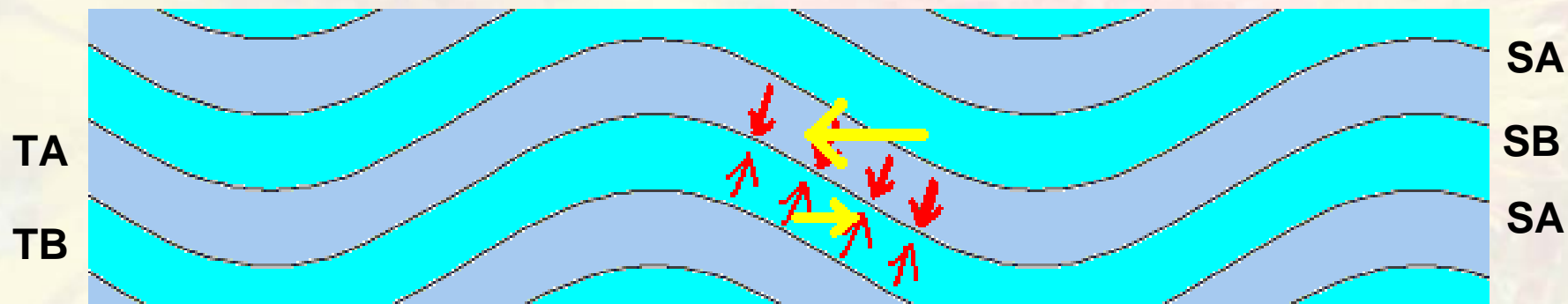
In-phase step wandering induced by the drift without evaporation

- Net diffusion current in x-direction

$$J_x^A + J_x^B \propto (D_{\perp} - D_{\parallel}) \Delta c \tan \theta \propto \frac{\partial y_s(x)}{\partial x}$$

- Restoring current due to the step stiffness

$$j_x \propto -\frac{\partial \mu}{\partial x} \propto -\frac{\partial}{\partial x} \left(\tilde{\beta} \frac{\partial^2 y_s}{\partial x^2} \right)$$



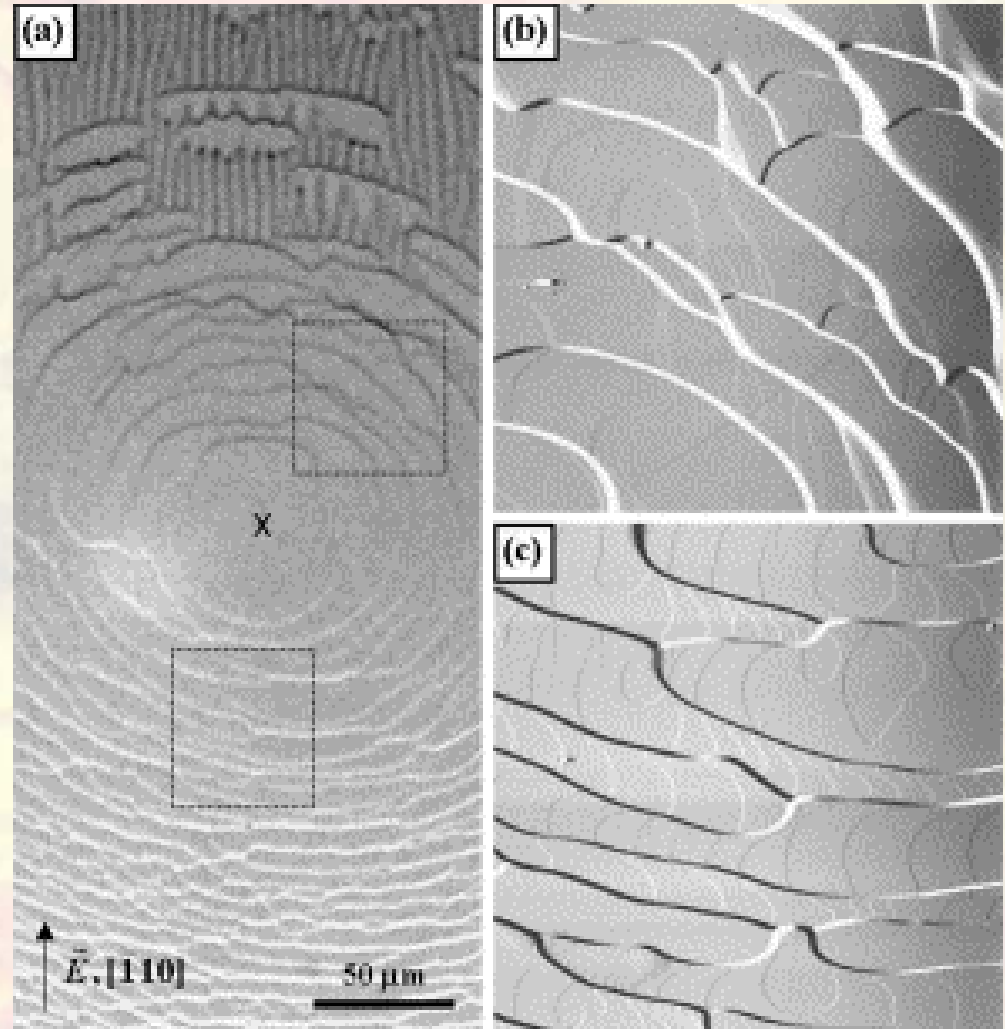
Wandering motion by the transverse currents

- The balance of adatom currents in x -direction

$$\begin{aligned}
 \frac{\partial y_s}{\partial t} &\propto -\nabla \cdot J \\
 &= -\frac{\partial}{\partial x} (J_x^A + J_x^B + J_x^{relax}) \\
 &\propto -\frac{\partial}{\partial x} \left(\frac{\partial y_s(x)}{\partial x} - \frac{\partial}{\partial x} \left(\frac{\partial^2 y_s(x)}{\partial x^2} \right) \right) \\
 &= -\frac{\partial^2 y_s(x)}{\partial x^2} + \frac{\partial^4 y_s(x)}{\partial x^4} \\
 &\propto (k^2 - k^4) y_s(k)
 \end{aligned}$$

Si(001) surface after CD heating

- Dimpled Si(0 0 1) surface heated to 990 ° C for 18 h with the current along $[-110]$.
- (b) and (c) "Illuminated" from the left.



Nonlinear effect in conservative systems

- Velocity of a step

$$\frac{\partial y_s(x)}{\partial t} = -\frac{\partial}{\partial x} J_x = -\frac{\partial}{\partial x} \left(J_x^{growth} + J_x^{relax} \right)$$

- Full equation (after scale transformation)

$$\frac{\partial y_s}{\partial t} = -\frac{\partial}{\partial x} \left[\frac{1}{1 + \left(\frac{\partial y_s}{\partial x} \right)^2} \left(\frac{\partial y_s}{\partial x} + \frac{\partial}{\partial x} \left(\frac{\frac{\partial^2 y_s}{\partial x^2}}{\left(1 + \left(\frac{\partial y_s}{\partial x} \right)^2 \right)^{3/2}} \right) \right) \right]$$

Instabilities by drift in MC simulation with strong step repulsion

Step down

Step up

Heavy figures deleted See

M. Sato, M. Uwaha, Y. Saito and Y. Hirose, Phys. Rev. B 67, 125408 (2003)

Evolution of wandering pattern---coarsening

$t = 2.5 \times 10^5$

$t = 12.5 \times 10^5$

Heavy figures deleted See

M. Sato, M. Uwaha, Y. Saito and Y. Hirose, Phys. Rev. B 67, 125408 (2003)

Evolution of wandering ---coarsening exponents

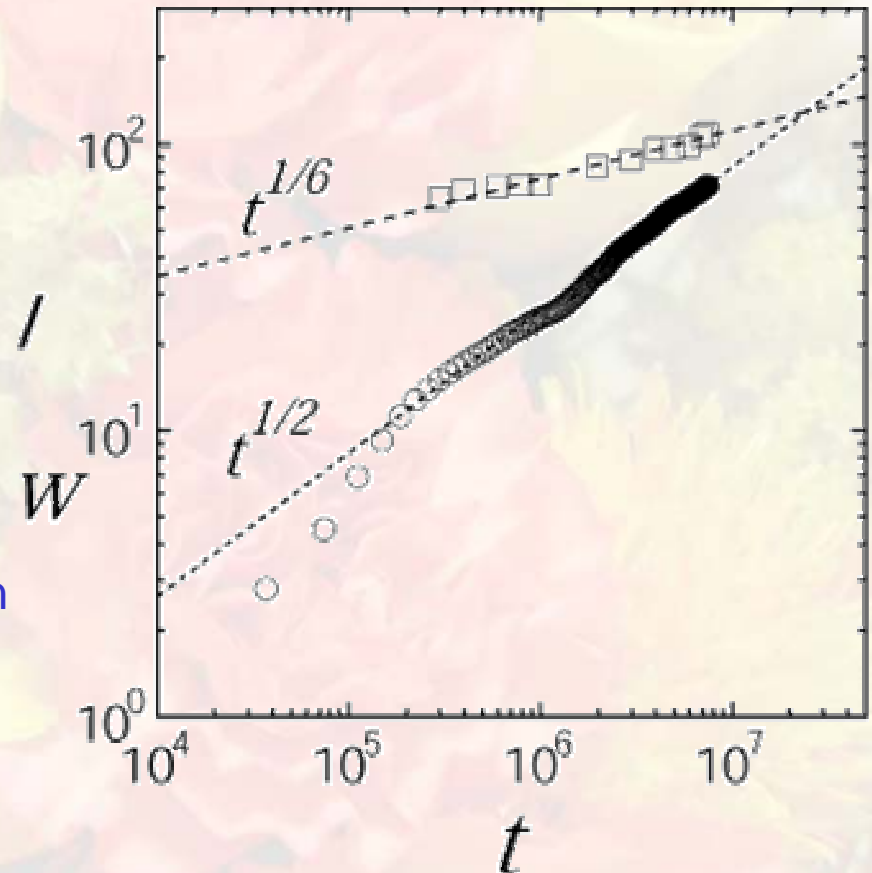
Wandering amplitude

$$W \propto t^{1/2}$$

Wandering wavelength

$$\lambda \propto t^{1/6}$$

Agree with the result of numerical integration of the continuum equation (Paulin, Gillet, Pierre-Louis, Misbah, Phys. Rev. Lett. 86 (2001) 5538)



Summary

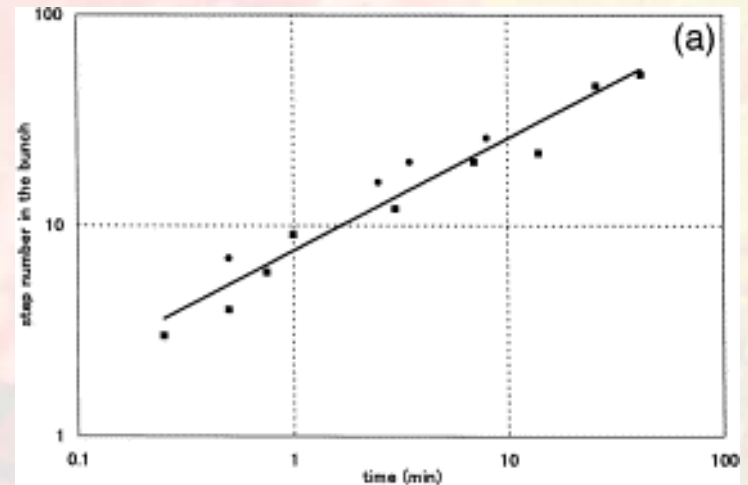
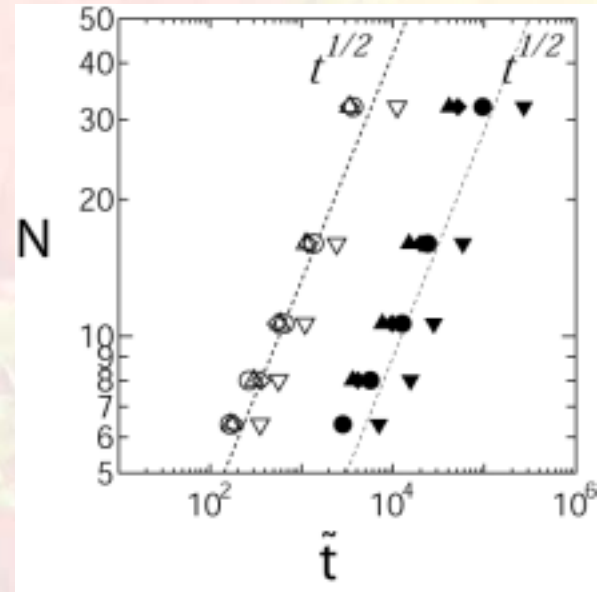
- Formation of large bunches in **both** drift directions is explained by a simple picture.
- Under strong step repulsion, bunching is suppressed but **in-phase step wandering** occurs with step-up drift
Periodic growing pattern appears.

**Imbalance of diffusion current by the step repulsion
(many-body effect!)**

(in agreement with experiment)

Remaining problem

- Why is the time scale for bunch growth independent of the drift direction?





Thanks to CSCAMM