<u>Well-Balanced Positivity Preserving Central-</u> <u>Upwind Scheme on Triangular Grids for the</u> <u>Saint-Venant System</u>

Yekaterina Epshteyn, University of Utah joint work with Steve Bryson, Alexander Kurganov and Guergana Petrova

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<u>Outline</u>

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Motivation

- Saint-Venant System of shallow water equations describes the fluid flow as a conservation law with an additional source term
- The general characteristic of shallow water flows is that vertical scales of motion are much smaller than the horizontal scales
- The shallow water equations are derived from the incompressible Navier-Stokes

<u>Motivation</u>

- This Saint-Venant System is widely used in many scientific and engineering applications related to
- Modeling of water flows in rivers, lakes and coastal areas
- The Development of robust and accurate numerical methods for Shallow Water Equations is an important and challenging problem



<u>2-D Saint-Venant system of shallow water</u> <u>equations</u>

$$\begin{cases} h_t + (hu)_x + (hv)_y = 0, \\ (hu)_t + \left(hu^2 + \frac{1}{2}gh^2\right)_x + (huv)_y = -ghB_x, \\ (hv)_t + (huv)_x + \left(hv^2 + \frac{1}{2}gh^2\right)_y = -ghB_y, \end{cases}$$
(1)

- \bullet the function B(x,y) represents the bottom elevation
- $\bullet~h$ is the fluid depth above the bottom
- $(u, v)^T$ is the velocity vector
- $\bullet~g$ is the gravitational constant

One of the difficulties encountered:

- that system (1) admits nonsmooth solutions: shocks, rarefaction waves,
- the bottom topography function *B* can be discontinuous.

2-D Saint-Venant system of shallow water

equations

A good numerical method for Saint-Venant System should have at least two major properties, which are crucial for its stability:

(i) The method should be well-balanced, that is, it should exactly preserve the stationary steady-state solutions $h + B \equiv \text{const}, \ u \equiv v \equiv 0$ (lake at rest states).

This property diminishes the appearance of unphysical waves of magnitude proportional to the grid size (the so-called "numerical storm"), which are normally present when computing quasi steady-states;

(ii) The method should be positivity preserving, that is, the water depth h should be nonnegative at all times.

This property ensures a robust performance of the method on dry (h = 0) and almost dry $(h \sim 0)$ states.

<u>Semi-discrete central-upwind scheme</u>

Central-Upwind schemes were developed for multidimensional hyperbolic systems of conservation laws in 2000 - 2007 by Kurganov, Lin, Noelle, Petrova, Tadmor, ...

- Central-Upwind schemes are Godunov-type finitevolume projection-evolution methods:
- At each time level a solution is globally approximated by a piecewise polynomial function,
- Which is then evolved to the new time level using the integral form of the conservation law system.

<u>Key ideas of the scheme development for</u> <u>Saint-Venant system</u>

- Change of conservative variables from $(h, hu, hv)^T$ to $(w := h + B, hu, hv)^T$
- Replacement of the bottom topography function *B* with its continuous piecewise linear (or bilinear in the 2-D case) approximation
- Special positivity preserving correction of the piecewise linear reconstruction for the water surface w
- Development of a special finite-volume-type quadrature for the discretization of the cell averages of the geometric source term.

Description of the scheme

- We describe now, our new second-order semidiscrete central-upwind scheme for solving the Saint-Venant system of shallow water equations on triangular grids
- We first denote the water surface by w := h + B and rewrite the original Saint-Venant system in terms of the vector $\mathbf{U} := (w, hu, hv)^T$:

$$\mathbf{U}_t + \mathbf{F}(\mathbf{U}, B)_x + \mathbf{G}(\mathbf{U}, B)_y = \mathbf{S}(\mathbf{U}, B)$$

where the fluxes and the source terms are:

$$\mathbf{F}(\mathbf{U}, B) = \left(hu, \frac{(hu)^2}{w - B} + \frac{1}{2}g(w - B)^2, \frac{(hu)(hv)}{w - B}\right)^T$$
$$\mathbf{G}(\mathbf{U}, B) = \left(hv, \frac{(hu)(hv)}{w - B}, \frac{(hv)^2}{w - B} + \frac{1}{2}g(w - B)^2\right)^T$$
$$\mathbf{S}(\mathbf{U}, B) = \left(0, -g(w - B)B_x, -g(w - B)B_y\right)^T.$$

Description of the scheme: notations



- $\vec{n}_{jk} := (\cos(\theta_{jk}), \sin(\theta_{jk}))$ are the outer unit normals to the corresponding sides of T_j of length ℓ_{jk} , k = 1, 2, 3,
- (x_j, y_j) are the coordinates of the center of mass for T_j and $M_{jk} = (x_{jk}, y_{jk})$ is the midpoint of the *k*-th side of the triangle T_j , k = 1, 2, 3
- T_{j1} , T_{j2} and T_{j3} are the neighboring triangles that share a common side with T_j

<u>Description of the central-upwind scheme on</u> <u>triangular grids</u>

Denote $\overline{\mathbf{U}}_{j}(t) \approx \frac{1}{|T_{j}|} \int_{T_{j}} \mathbf{U}(x, y, t) \, dx \, dy.$

Second order central-upwind scheme on triangular grid for the Saint-Venant System:

$$\frac{d\overline{\mathbf{U}}_j}{dt} =$$

$$-\frac{1}{|T_{j}|}\sum_{k=1}^{3}\frac{\ell_{jk}\cos(\theta_{jk})}{a_{jk}^{\text{in}}+a_{jk}^{\text{out}}}\left[a_{jk}^{\text{in}}\mathbf{F}(\mathbf{U}_{jk}(M_{jk}),B(M_{jk}))+a_{jk}^{\text{out}}\mathbf{F}(\mathbf{U}_{j}(M_{jk}),B(M_{jk}))\right] \\ -\frac{1}{|T_{j}|}\sum_{k=1}^{3}\frac{\ell_{jk}\sin(\theta_{jk})}{a_{jk}^{\text{in}}+a_{jk}^{\text{out}}}\left[a_{jk}^{\text{in}}\mathbf{G}(\mathbf{U}_{jk}(M_{jk}),B(M_{jk}))+a_{jk}^{\text{out}}\mathbf{G}(\mathbf{U}_{j}(M_{jk}),B(M_{jk}))\right] \\ +\frac{1}{|T_{j}|}\sum_{k=1}^{3}\ell_{jk}\frac{a_{jk}^{\text{in}}a_{jk}^{\text{out}}}{a_{jk}^{\text{in}}+a_{jk}^{\text{out}}}\left[\mathbf{U}_{jk}(M_{jk})-\mathbf{U}_{j}(M_{jk})\right] + \overline{\mathbf{S}}_{j},$$

<u>Description of the central-upwind scheme on</u> <u>triangular grids</u>

• $U_j(M_{jk})$ and $U_{jk}(M_{jk})$ are the corresponding values at M_{jk} of the piecewise linear reconstruction

 $\widetilde{\mathbf{U}}(x,y) := \overline{\mathbf{U}}_j + (\mathbf{U}_x)_j (x - x_j) + (\mathbf{U}_y)_j (y - y_j), \quad (x,y) \in T_j$ of U at time t

- The quantity $\overline{\mathbf{S}}_j$ in the scheme is an appropriate discretization of the cell averages of the source term
- The directional local speeds a_{jk}^{in} and a_{jk}^{out} are defined by

 $a_{jk}^{\rm in}(M_{jk}) = -\min\{\lambda_1[V_{jk}(\mathbf{U}_j(M_{jk}))], \lambda_1[V_{jk}(\mathbf{U}_{jk}(M_{jk})], 0\},$

 $a_{jk}^{\text{out}}(M_{jk}) = \max\{\lambda_3[V_{jk}(\mathbf{U}_j(M_{jk}))], \lambda_3[V_{jk}(\mathbf{U}_{jk}(M_{jk})], 0\},$

where $\lambda_1 [V_{jk}] \leq \lambda_2 [V_{jk}] \leq \lambda_3 [V_{jk}]$ are the eigenvalues of the matrix $V_{jk} = \cos(\theta_{jk}) \frac{\partial \mathbf{F}}{\partial \mathbf{U}} + \sin(\theta_{jk}) \frac{\partial \mathbf{G}}{\partial \mathbf{U}}$.

• A fully discrete scheme is obtained by using a stable ODE solver of an appropriate order

<u>Calculation of the numerical derivatives of the ith</u> <u>component of U</u>

- Construct three linear interpolations $L_j^{12}(x,y)$, $L_j^{23}(x,y)$ and $L_j^{13}(x,y)$: conservative on T_j and two of the neighboring triangles (T_{j1}, T_{j2}) , (T_{j2}, T_{j3}) and (T_{j1}, T_{j3})
- Select the linear piece with the smallest magnitude of the gradient, say, $L_i^{km}(x, y)$, and set

 $((\mathbf{U}_x^{(i)})_j, (\mathbf{U}_y^{(i)})_j)^T = \nabla L_j^{km}$

• Minimize the oscillations by checking the appearance of local extrema at the points M_{jk} , 1,2,3

Piecewise linear approximation of the bottom

• Replace the bottom topography function B with its continuous piecewise linear approximation \tilde{B} , which over each cell T_i is given by the formula:

$$\begin{array}{c|cccc} x - \tilde{x}_{j_{12}} & y - \tilde{y}_{j_{12}} & \widetilde{B}(x, y) - \mathcal{B}_{j_{12}} \\ \tilde{x}_{j_{23}} - \tilde{x}_{j_{12}} & \tilde{y}_{j_{23}} - \tilde{y}_{j_{12}} & \mathcal{B}_{j_{23}} - \mathcal{B}_{j_{12}} \\ \tilde{x}_{j_{13}} - \tilde{x}_{j_{12}} & \tilde{y}_{j_{13}} - \tilde{y}_{j_{12}} & \mathcal{B}_{j_{13}} - \mathcal{B}_{j_{12}} \\ \end{array} \right| = 0, \quad (x, y) \in T_j.$$

- $\mathcal{B}_{j_{\kappa}}$ are the values of \widetilde{B} at the vertices $(\tilde{x}_{j_{\kappa}}, \tilde{y}_{j_{\kappa}}), \kappa = 12, 23, 13$, of the cell T_j
- $\mathcal{B}_{j_{\kappa}} := \frac{1}{2} (\max_{\xi^2 + \eta^2 = 1} \lim_{h, \ell \to 0} B(\tilde{x}_{j_{\kappa}} + h\xi, \tilde{y}_{j_{\kappa}} + \ell\eta) + \min_{\xi^2 + \eta^2 = 1} \lim_{h, \ell \to 0} B(\tilde{x}_{j_{\kappa}} + h\xi, \tilde{y}_{j_{\kappa}} + \ell\eta)),$
- If the function B is continuous at $(\tilde{x}_{j_{\kappa}}, \tilde{y}_{j_{\kappa}})$: $\mathcal{B}_{j_{\kappa}} = B(\tilde{x}_{j_{\kappa}}, \tilde{y}_{j_{\kappa}})$

Positivity preserving reconstruction for w

The idea of the algorithm that guarantees positivity of the reconstructed values of the water depth $h_j(M_{jk}) := w_j(M_{jk}) - B_{jk}, \ k = 1, 2, 3$, for all *j*:

- The reconstruction \widetilde{w} should be corrected only in those triangles, where $\widetilde{w}(\widetilde{x}_{j_{\kappa}}, \widetilde{y}_{j_{\kappa}}) < \mathcal{B}_{j_{\kappa}}$ for some κ , $\kappa = 12, 23, 13$
- Since $\overline{w}_j \ge B_j$, it is impossible to have $\widetilde{w}(\widetilde{x}_{j_{\kappa}}, \widetilde{y}_{j_{\kappa}}) < \mathcal{B}_{j_{\kappa}}$ for all three values of κ : at all three vertices of the triangle T_j
- Two cases in which a correction is needed are possible:

either there are two indices κ_1 and κ_2 , for which $\widetilde{w}(\widetilde{x}_{j_{\kappa_1}}, \widetilde{y}_{j_{\kappa_1}}) < \mathcal{B}_{j_{\kappa_1}}$ and $\widetilde{w}(\widetilde{x}_{j_{\kappa_2}}, \widetilde{y}_{j_{\kappa_2}}) < \mathcal{B}_{j_{\kappa_2}}$, or there is only one index κ_1 , for which $\widetilde{w}(\widetilde{x}_{j_{\kappa_1}}, \widetilde{y}_{j_{\kappa_1}}) < \mathcal{B}_{j_{\kappa_1}}$

Well-balanced discretization of the source term

- The well-balanced property of the scheme is guaranteed if the discretized cell average of the source term, $\overline{\mathbf{S}}_j$, exactly balances the numerical fluxes
- The desired quadrature for the source term that will preserve stationary steady states $(\mathbf{U}_{jk}(M_{jk}) \equiv \mathbf{U}_j(M_{jk}) \equiv (C, 0, 0)^T, \forall j, k)$ is given by:

$$\overline{\mathbf{S}}_{j}^{(2)} = \frac{g}{2|T_{j}|} \sum_{k=1}^{3} \ell_{jk} (w_{j}(M_{jk}) - B_{jk})^{2} \cos(\theta_{jk}) - g(w_{x})_{j} (\overline{w}_{j} - B_{j})$$
$$\overline{\mathbf{S}}_{j}^{(3)} = \frac{g}{2|T_{j}|} \sum_{k=1}^{3} \ell_{jk} (w_{j}(M_{jk}) - B_{jk})^{2} \sin(\theta_{jk}) - g(w_{y})_{j} (\overline{w}_{j} - B_{j})$$

Main theorem: positivity property of the new scheme

Theorem 1 Consider the Saint-Venant system in the new variables $\mathbf{U} := (w, hu, hv)^T$ and the central-upwind semi-discrete scheme (with well-balanced quadrature for the source S, positivity preserving reconstruction for w)

- Assume that the system of ODEs for the fully discrete scheme is solved by the forward Euler method and that for all j, $\overline{w}_j^n - B_j \ge 0$ at time $t = t^n$
- Then, for all j, $\overline{w}_{j}^{n+1} B_{j} \ge 0$ at time $t = t^{n+1} = t^{n} + dt$, provided that $dt \le \frac{1}{6a} \min_{j,k} \{r_{jk}\}$, where $a := \max_{j,k} \{a_{jk}^{\text{out}}, a_{jk}^{\text{in}}\}$ and r_{jk} , k = 1, 2, 3, are the altitudes of triangle T_{j}

Remark. Theorem 1 is still valid if one uses a higher-order SSP ODE solver (either the Runge-Kutta or the multistep one), because such solvers can be written as a convex combination of several forward Euler steps.

Accuracy test

The scheme is applied to the Saint-Venant system subject to the following initial data and the bottom topography:

$$w(x, y, 0) = 1, \quad u(x, y, 0) = 0.3,$$

 $B(x, y) = 0.5 \exp(-25(x - 1)^2 - 50(y - 0.5)^2).$

- For a reference solution, we solve this problem with our method on a $2 \times 400 \times 400$ triangular grid
- By t = 0.07 the solution converges to the steady state

Accuracy test

• w component of the reference solution of the IVP on a $2 \times 400 \times 400$ grid: the 3-D view (left) and the contour plot (right).



• L^1 - and L^∞ -errors and numerical orders of accuracy.

Number of cells	L^1 -error	Order	L^{∞} -error	Order
$2 \times 50 \times 50$	6.59e-04	_	8.02e-03	_
$2 \times 100 \times 100$	2.87e-04	1.20	3.59e-03	1.16
$2 \times 200 \times 200$	1.00e-04	1.52	1.21e-03	1.57

<u>Small perturbation of a stationary steady-state</u> <u>solution</u>

- Solve the initial value problem (IVP) proposed by R.Leveque.
- The computational domain is $[0,2] \times [0,1]$ and the bottom consists of an elliptical shaped hump:

$$B(x, y) = 0.8 \exp(-5(x - 0.9)^2 - 50(y - 0.5)^2).$$

• Initially, the water is at rest and its surface is flat everywhere except for 0.05 < x < 0.15:

 $w(x, y, 0) = \begin{cases} 1 + \varepsilon, & 0.05 < x < 0.15, \\ 1, & \text{otherwise,} \end{cases} \quad u(x, y, 0) \equiv v(x, y, 0) \equiv 0, \end{cases}$

where the perturbation height is $\varepsilon = 10^{-4}$

<u>Perturbation of a stationary steady-state: well-balanced</u> <u>scheme (left) and non well-balanced (right)</u>











<u>Perturbation of a stationary steady-state: well-</u> balanced scheme (left) and non well-balanced (right)









<u>Saint-Venant System with friction and</u> <u>discontinuous bottom</u>

- More realistic shallow water models include additional friction and/or viscosity terms
- Presence of friction and viscosity terms guarantees uniqueness of the steady state solution
- We consider the simplest model in which only friction terms, $-\kappa(h)u$ and $-\kappa(h)v$, are added to the rhs of the second and third equations of the Saint-Venant System

$$\begin{cases} h_t + (hu)_x + (hv)_y = 0, \\ (hu)_t + \left(hu^2 + \frac{1}{2}gh^2\right)_x + (huv)_y = -ghB_x - \kappa(h)u, \\ (hv)_t + (huv)_x + \left(hv^2 + \frac{1}{2}gh^2\right)_y = -ghB_y - \kappa(h)v. \end{cases}$$

Saint-Venant System with friction and discontinuous

<u>bottom</u>

- We numerically solve the shallow water model with friction term on the domain $[-0.25, 1.75] \times [-0.5, 0.5]$
- We assume that the friction coefficient is

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\kappa(h) = 0.001(1+10h)^{-1}
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• The bottom topography function has a discontinuity along the vertical line x = 1 and it mimics a mountain river valley





<u>Saint-Venant System with friction and discontinuous</u> <u>bottom: description of the initial and boundary</u> <u>conditions</u>

- We implement reflecting (solid wall) boundary conditions at all boundaries
- Our initial data correspond to the situation when the second of the three dams, initially located at the vertical lines

x = -0.25 (the left boundary of the computational domain), x = 0, and x = 1.75 (the right boundary of the computational domain),

breaks down at time t = 0, and the water propagates into the initially dry area x > 0, and a "lake at rest" steady state is achieved after a certain period of time

- We plot 1-D slices of the numerical solution along the y = 0 line
- Plots clearly show the dynamics of the fluid flow as it moves from the region x < 0 into the initially dry area x > 0 and gradually settles down into a "lake at rest" steady state

— t=0

t=1

t=2

- t=4

1.6

1.6

1.6

1.6

steady state

1.6

1.2

1.2

1.2

12

1.4

1.4

1.4

1.4

1.4



• This state includes dry areas and therefore its computation requires a method that is both well-balanced and positivity preserving on the entire computational domain

Flow in converging-diverging channel

• The exact geometry of each channel is determined by its breadth, which is equal to $2y_b(x)$, where

$$y_{\rm b}(x) = \begin{cases} 0.5 - 0.5(1 - d)\cos^2(\pi(x - 1.5)), & |x - 1.5| \le 0.5, \\ 0.5, & \text{otherwise}, \end{cases}$$

• d = 0.6 is the minimum channel breadth





Flow in converging-diverging channel

• The initial conditions:

 $w(x, y, 0) = \max\left\{1, B(x, y)\right\}, \quad u(x, y, 0) = 2, \quad v(x, y, 0) = 0.$

- The upper and lower y-boundaries are reflecting (solid wall), the left x-boundary is an inflow boundary with u = 2 and the right x-boundary is a zero-order outflow boundary
- The bottom topography is given by

$$B(x,y) = \left(e^{-10(x-1.9)^2 - 50(y-0.2)^2} + e^{-20(x-2.2)^2 - 50(y+0.2)^2}\right),$$



Flow in converging-diverging channel: w

Steady-state solution (w) for $(d, B_{\text{max}}) = (0.6, 1)$ on $2 \times 200 \times 200$ (left) and $2 \times 400 \times 400$ (right) grids.



<u>Conclusions/Difficulties</u>

- We developed a simple central-upwind scheme for the Saint-Venant system on triangular grids
- We proved that the scheme both preserves stationary steady states (lake at rest) and guarantees the positivity of the computed fluid depth
- It can be applied to models with discontinuous bottom topography and irregular channel widths
- Method is sensitive to the accuracy of the boundary representation
- S. Bryson, Y. Epshteyn, A. Kurganov and G. Petrova, Well-Balanced Positivity Preserving Central-Upwind Scheme on Triangular Grids for the Saint-Venant System, to appear, ESAIM: M2AN 2010.