#### **Introduction of Particle Image Velocimetry**



#### Burgers Program For Fluid Dynamics Turbulence School College Park, Maryland, May 24-27

Slides largely generated by J. Westerweel & C. Poelma of Technical University of Delft Adapted by K. Kiger



#### Particle Image Velocimetry (PIV):

Imaging of tracer particles, calculate displacement: local fluid velocity



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#### Particle Image Velocimetry (PIV)





- divide image pair in *interrogation regions*
- small region:
   uniform motion
- compute displacement
- repeat !!!

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Particle Image Velocimetry (PIV):



Instantaneous measurement of 2 components in a plane

# **conventional methods** (HWA, LDV)

- single-point measurement
- traversing of flow domain
- time consuming
- only turbulence statistics

#### particle image velocimetry

- whole-field method
- non-intrusive (seeding
- instantaneous flow field





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Particle Image Velocimetry (PIV):

Instantaneous measurement of 2 components in a plane



#### particle image velocimetry

- whole-field method
- non-intrusive (seeding
- instantaneous flow field



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PIV components:

- tracer particles
- light source
- light sheet optics
- camera
- measurement settings
- interrogation
- post-processing



Hardware (imaging)

Software (image analysis)

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### **Tracer particles**

#### **Assumptions:**

- homogeneously distributed
- follow flow perfectly
- uniform displacement within interrogation region

#### Criteria:

-easily visible-particles should not influence fluid flow!

small, volume fraction <  $10^{-4}$ 

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# **Evaluation at higher density**

High N<sub>I</sub>: no longer possible/desirable to follow individual tracer particles



Particle can be matched with a number of candidates

Repeat process for other particles, sum up: "wrong" combinations will lead to noise, but "true" displacement will dominate

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### Statistical estimate of particle motion

 Statistical correlations used to find average particle displacement

$$R(i,j) = \frac{\sum_{k=1}^{B_x} \sum_{l=1}^{B_y} (I_a(k,l) - \bar{I}_a) (I_b(k+i,l+j) - \bar{I}_b)}{\left[\sum_{k=1}^{B_x} \sum_{l=1}^{B_y} (I_a(k,l) - \bar{I}_a)^2 \sum_{k=1}^{B_x} \sum_{l=1}^{B_y} (I_b(k+i,l+j) - \bar{I}_b)^2\right]^{\frac{1}{2}}}$$

$$\bar{I}_{a} = \frac{1}{B_{x}B_{y}} \sum_{k=1}^{B_{x}} \sum_{l=1}^{B_{y}} I_{a}(k,l)$$

1-d image @ t=t<sub>0</sub> 300 Amplitude (A.U.) 00 00 00 -20 -1020 30 -3010 x position (pixels) 1-d image @ t=t1 300 Amplitude (A.U.) -20 -10 10 20 30 0 x position (pixels) **Cross-correlation** Cross-correlation Coefficient, R( Δx) 0.8 0.6 0.4 0.2 -0.2 -0.4 -20 -15 10 -10 0 15 20 -5 5

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Correlation displacement, Ax



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#### **1-D cross-correlation example**



### Finding the maximum displacement

-Shift 2<sup>nd</sup> window with respect to the first

Typically 16x16 or 32x32 pixels

- Calculate "match"

Good indicator: R(i,j)

$$P) = \frac{\sum_{k=1}^{N} \sum_{l=1}^{n} (I_a(k,l) - \bar{I}_a) (I_b(k+i,l+j) - \bar{I}_b)}{\left[\sum_{k=1}^{B_x} \sum_{l=1}^{B_y} (I_a(k,l) - \bar{I}_a)^2 \sum_{k=1}^{B_x} \sum_{l=1}^{B_y} (I_b(k+i,l+j) - \bar{I}_b)^2\right]^{\frac{1}{2}}}$$

- Repeat to find best estimate

$$\bar{I}_{a} = \frac{1}{B_{x}B_{y}} \sum_{k=1}^{B_{x}} \sum_{l=1}^{B_{y}} I_{a}(k,l)$$

 $B_x \quad B_y$ 



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### Finding the maximum displacement

-Shift 2<sup>nd</sup> window with respect to the first

Typically 16x16 or 32x32 pixels

- Calculate "match"

Good indicator:

$$R(i,j) = \frac{\sum_{k=1}^{n} \sum_{l=1}^{m} (I_a(k,l) - \bar{I}_a) (I_b(k+i,l+j) - \bar{I}_b)}{\left[\sum_{k=1}^{B_x} \sum_{l=1}^{B_y} (I_a(k,l) - \bar{I}_a)^2 \sum_{k=1}^{B_x} \sum_{l=1}^{B_y} (I_b(k+i,l+j) - \bar{I}_b)^2\right]^{\frac{1}{2}}}$$

- Repeat to find best estimate



$$\bar{I}_{a} = \frac{1}{B_{x}B_{y}} \sum_{k=1}^{B_{x}} \sum_{l=1}^{B_{y}} I_{a}(k,l)$$

Bad match: sum of product of intensities low

 $B_r B_v$ 

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## Finding the maximum displacement

-Shift 2<sup>nd</sup> window with respect to the first

Typically 16x16 or 32x32 pixels

- Calculate "match"

Good indicator: R(i,j)

$$=\frac{\sum_{k=1}^{B_{x}} \left(I_{a}(k,l) - \bar{I}_{a}\right) \left(I_{b}(k+i,l+j) - \bar{I}_{b}\right)}{\left[\sum_{k=1}^{B_{x}} \sum_{l=1}^{B_{y}} \left(I_{a}(k,l) - \bar{I}_{a}\right)^{2} \sum_{k=1}^{B_{x}} \sum_{l=1}^{B_{y}} \left(I_{b}(k+i,l+j) - \bar{I}_{b}\right)^{2}\right]^{\frac{1}{2}}}$$

- Repeat to find best estimate

$$\bar{I}_{a} = \frac{1}{B_{x}B_{y}} \sum_{k=1}^{B_{x}} \sum_{l=1}^{B_{y}} I_{a}(k,l)$$



Good match: sum of product of intensities high

 $B_x \quad B_y$ 

#### ◆Can be implemented as 2D FFT for digitized data

 Impose periodic conditions on interrogation region...causes bias error if not treated properly.

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#### **Cross-correlation**



This "shifting" method can formally be expressed as a cross-correlation:

$$R(\mathbf{s}) = \int I_1 \mathbf{x} I_2 \mathbf{x} + \mathbf{s} d\mathbf{x}$$

- $I_1$  and  $I_2$  are interrogation areas (sub-windows) of the total frames
- x is interrogation location
- s is the shift between the images

"Backbone" of PIV:

-cross-correlation of interrogation areas -find location of displacement peak

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particle concentration light sheet thickness int. area size magnification



More particles: better signal-to-noise ratio

Unambiguous detection of peak from noise:  $N_I=10$  (average), minimum of 4 per area in 95% of areas (number of tracer particles is a Poisson distribution)

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PTV: 1 particle used for velocity estimate; error *e* PIV: error ~ *e*/sqrt(N<sub>I</sub>)

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| $\Delta X / D_I = 0.00$    | 0.28  | 0.56  | 0.85  |
|----------------------------|---|---|---|
| $F_{I} = 1.00$             | 0.64  | 0.36  | 0.16  |
| $R_D(\mathbf{s}_D) \sim N$ | $V_I F_I \Longrightarrow F_I (\Delta X, \Delta$ | $(X) = \left(1 - \frac{ \Delta X }{D_I}\right) \left(1 - \frac{ \Delta X }{D_I}\right)$ | $\left. \frac{\Delta Y \Big }{D_I} \right)$ |

### Influence of out-of-plane displacement



Z-Displacement

<

quarter of light sheet thickness  $(\Delta z_0)$ 



 $\Delta Z / \Delta z_0 = 0.00 \qquad 0.25 \qquad 0.50 \qquad 0.75$  $F_0 = 1.00 \qquad 0.75 \qquad 0.50 \qquad 0.25$  $R_D(\mathbf{s}_D) \sim N_I F_I F_0 \implies F_0(\Delta z) = 1 - \frac{|\Delta z|}{\Delta z_0}$ 

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#### **PIV "Design rules"**



- image density •
- in-plane motion
- out-of-plane motion  $|\Delta z| < \frac{1}{4} \Delta z_0$
- spatial gradients

*N*<sup>*i*</sup> >10

- $|\Delta X| < \frac{1}{4} D_1$
- $M_0 |\Delta u| \Delta t < d_\tau$

Obtained by Keane & Adrian (1993) using synthetic data

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### **Window shifting**



in-plane motion

# strongly limits dynamic range of PIV



large window size: too much spatial averaging

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### **Window shifting**



in-plane motion

# $|\Delta X| < \frac{1}{4} D_{I}$

strongly limits dynamic range of PIV



small window size: too much in-plane pair loss

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### **Window shifting**

< <sup>1</sup>/<sub>4</sub> D<sub>1</sub>

in-plane motion

# $|\Delta X| < \frac{1}{4} D_{l}$

# strongly limits dynamic range of PIV



Multi-pass approach: start with large windows, use this result as "pre-shift' for smaller windows...

No more in-plane pair loss limitations!

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### Window shifting: Example



Grid turbulence

#### fixed windows



windows at same location

#### matched windows



windows at 7px "downstream'

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#### Window shifting: Example



Vortex ring, decreasing window sizes



Raffel, Willert and Kompenhans

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### **Sub-pixel accuracy**



Maximum in the correlation plane: single-pixel resolution of displacement?



But the peak contains a lot more information!

Gaussian particle images  $\rightarrow$  Gaussian correlation peak (but smeared)

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Fractional displacement can be obtained using the distribution of gray values around maximum

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three-point estimators

peak centroid

$$arepsilon = rac{R_{+1} - R_{-1}}{R_{-1} + R_0 + R_{+1}}$$

parabolic peak fit

$$\varepsilon = \frac{R_{-1} - R_{+1}}{2 \, \P_{-1} + R_{+1} - 2R_0}$$

Gaussian peak fit

$$\varepsilon = \frac{\ln R_{-1} - \ln R_{+1}}{2 \ln R_{-1} + \ln R_{+1} - 2 \ln R_{0}}$$

$$\varepsilon \propto \frac{\text{balance}}{\text{normalization}}$$

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#### **Peak locking**



"zig-zag" structure, sudden "kinks" in the flow THE A. JAMES CLARK SCHOOL of ENGINEERING

# **Sub-pixel Interpolation Errors**

#### Accuracy depends on:

- particle image size
- noise in data (seeding density, camera noise)
- shear rate

#### Can exhibit "peak locking"

- Interpolation of peak is biased towards a symmetric data distribution (Integer and 1/2 integer peak locations)
- Polynomials exhibit strong locking when particle diameter is small
- Gaussian is most commonly used
- Splines are very robust, but expensive to calculate
- See Particle Image Velocimetry, by Raffel, Willert, and Kompenhans, Springer-Verlag, 1998.









#### 30000 30000 number of vectors number of vectors 20000 20000 10000 10000 0 0 10 12 14 4 8 12 10 14 displacement [px] displacement [px]

#### Histogram of velocities in a turbulent flow

centroid

Gaussian peak fit

Even with Gaussian peak fit:

particle image size too small  $\rightarrow$  peak locking (Consider a "point particle' sampled by discrete pixels)

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# Sub-pixel accuracy

optimal resolution: particle image size: ~2 px Smaller: particle no longer resolved Larger: random noise increase

"three-point" estimators:

Peak centroid Parabolic peak fit Gaussian peak fit



In practice 0.05-0.1 px

Main difference: sensitivity to "peak locking" or "pixel lock-in", bias towards integer displacements

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### application example: grid-generated turbulence



$$\Delta X = 7 \text{ px} \quad u'/U = 2.5\%$$

#### fixed windows





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### **Data Validation**

#### "article"





"lab"



Spurious or "Bad" vectors

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### **Spurious vectors**

Three main causes:

- insufficient particle-image pairs
- in-plane loss-of-pairs, out-of-plane loss-of-pairs
- gradients

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### **Effect of tracer density**



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#### **Remedies**



- increase  $N_{I}$ 
  - practical limitations:
    - optical transparency of the fluid
    - two-phase effects
    - image saturation / speckle
- detection, removal & replacement
  - keep finite  $N_l$  ( $\Gamma \sim 0.05$ )
    - data loss is small
    - signal loss occurs in isolated points
    - data recovery by interpolation

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#### **Detection methods**

- human perception
- peak height
  - amount of correlated signal
- peak detectability
  - peak height relative to noise
  - lower limit for SNR
- residual vector analysis
  - fluctuation of displacement
- multiplication of correlation planes
- fluid mechanics
  - continuity
- fuzzy logic & neural nets







### **Residual analysis**

- evaluate fluctuation of measured velocity  $\Rightarrow$  residual
- ideally:  $U_{ref}$  = true velocity
- reference values:
  - $U_{\rm ref}$  = global mean velocity
    - comparable to 2D-histogram analysis
    - · does not take local coherent motion into account
    - probably only works in homogeneous turbulence
  - U<sub>ref</sub> = local (3×3) mean velocity
    - · takes local coherent motion into account
    - · very sensitive to outliers in the local neighborhood
  - $U_{ref}$  = local (3×3) median velocity
    - · almost identical statistical properties as local mean
    - Strongly suppressed sensitivity to outliers in heighborhood



$$r = \left| U - U_{\text{ref}} \right|$$

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Standard mean and r.m.s. are very sensitive to bad data contamination... need robust measure of fluctuation

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#### **Median test**



- 1 Calculate reference velocity: median of 8 neighbors  $u_{ref} = median(u_1, u_2, ..., u_8)$
- 2 calculate residuals:  $r_i = u_i - u_{ref}$ 3 – Normalize target residual by: median(r<sub>i</sub>) +  $\varepsilon$

 $r_0^* = \frac{\left|u_0 - u_{ref}\right|}{median(r_i) + \varepsilon}$ 

4 – Robust measure found for:  $\varepsilon = 0.1$  and  $r_0^* > 2$ 



# Interpolation





Bilinear interpolation satisfies continuity

For 5% bad vectors, 80% of the vectors are isolated

Bad vector can be recovered without any problems

N.B.: interpolation biases statistics (power spectra, correlation function) Better not to replace bad vectors (use e.g. slotting method)

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# **Overlapping windows**



Method to increase data yield:

Allow overlap between adjacent interrogation areas



Motivation: particle pairs near edges contribute less to correlation result; Shift window so they are in the center: additional, relatively uncorrelated result

50% is very common, but beware of oversampling

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### A Generic PIV program





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#### **PIV software**

#### <u>Free</u>



PIVware: command line, linux (Westerweel) JPIV: Java version of PIVware (Vennemann) MatPiv: Matlab PIV toolbox (Cambridge, Sveen) URAPIV: Matlab PIV toolbox (Gurka and Liberzon) DigiFlow (Cambridge), PIV Sleuth (UIUC), MPIV, GPIV, CIV, OSIV,...

#### **Commercial**

PIVtec TSI Dantec LaVision Oxford Lasers/ILA PIVview Insight Flowmap DaVis VidPIV

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# **Particle Motion: tracer particle**

Equation of motion for spherical particle:



 Neglect: non-linear drag (only really needed for high-speed flows), Basset history term (higher order effect)

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# Simple thought experiment

- Lets see how a particle responds to a step change in velocity
  - Only consider viscous drag

$$m_p \frac{dv_p}{dt} = 3\pi\mu D(u - v_p) \qquad u = \begin{cases} 0 \text{ for } t < 0 \\ U \text{ for } t \ge 0 \end{cases}$$

$$\frac{dv_p}{dt} = \frac{18\mu}{\rho_p D^2} \left( U - v_p \right) = \frac{1}{\tau_p} \left( U - v_p \right)$$

$$v_p' + \frac{1}{\tau_p} v_p = \frac{1}{\tau_p} U$$

$$v_p = (v_p)_{\text{particular}} + (v_p)_{\text{homogeneous}} = U + C_1 \exp(-t/\tau_p)$$

$$v_p = U \left[ 1 - \exp\left(-t/\tau_p\right) \right]$$

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### **Particle Transfer Function**



- Useful to examine steady-state particle response to 1-D oscillating flow of arbitrary sum of frequencies
  - Represent *u* as an infinite sum of harmonic functions
  - Neglect gravity (DC response, not transient)
    - $u(t) = \int_{0}^{\infty} \Lambda_{f}(\omega) \exp(i\omega t) d\omega \qquad \qquad \frac{du(t)}{dt} = \int_{0}^{\infty} i\omega \Lambda_{f}(\omega) \exp(i\omega t) d\omega$  $v_{p}(t) = \int_{0}^{\infty} \Lambda_{p}(\omega) \exp(i\omega t) d\omega \qquad \qquad \frac{dv_{p}(t)}{dt} = \int_{0}^{\infty} i\omega \Lambda_{p}(\omega) \exp(i\omega t) d\omega$

$$\frac{2m_p}{3\pi\mu D}\frac{dv_p}{dt} = 2\left(u - v_p\right) + \frac{m_f}{m_p}\frac{m_p}{3\pi\mu D}\left[\frac{du}{dt} - \frac{dv_p}{dt}\right] + 2\frac{m_f}{m_p}\frac{m_p}{3\pi\mu D}\frac{du}{dt}$$

$$\int_{0}^{\infty} 2\frac{\rho_{p}D^{2}}{18\mu}\Lambda_{p}i\omega\exp(i\omega t)d\omega = \int_{0}^{\infty} 2(\Lambda_{f} - \Lambda_{p})\exp(i\omega t)d\omega + \int_{0}^{\infty} \frac{\rho_{f}}{\rho_{p}}\frac{\rho_{p}D^{2}}{18\mu}i\omega[3\Lambda_{f} - \Lambda_{p}]\exp(i\omega t)d\omega$$

$$2iSt\Lambda_{p} = 2(\Lambda_{f} - \Lambda_{p}) + i\gamma St[3\Lambda_{f} - \Lambda_{p}] \qquad St = \frac{\tau_{p}}{\tau_{f}} = \frac{\rho_{p}D^{2}\omega}{18\mu} \qquad \gamma = \frac{\rho_{f}}{\rho_{p}}$$

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### **Particle Transfer Function**



• This can be rearranged:

$$\left\| \frac{\mathbf{r}}{\mathbf{u}}_{p} \right\| = \left[ \frac{\Lambda_{p} \Lambda_{p}^{*}}{\Lambda_{f} \Lambda_{f}^{*}} \right]^{1/2} = \left[ \frac{A^{2} + B^{2}}{A^{2} + 1} \right]^{1/2} \qquad A = \frac{2}{St \ \mathbf{\ell} + \gamma}$$

$$\phi = \tan^{-1} \left[ \frac{\operatorname{Im}(\Lambda_{p} / \Lambda_{f})}{\operatorname{Re}(\Lambda_{p} / \Lambda_{f})} \right] = \tan^{-1} \left[ \frac{A(1-B)}{A^{2} + B} \right] \qquad B = \frac{3\gamma}{2 + \lambda}$$

#### **Examine**

- Liquid in air, gas in water, plastic in water

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### Liquid particles in Air



- Liquid drops in air require St~0.3 for 95% fidelity (1  $\mu$ m ~ 15 kHz)
- Size relatively unimportant for near-neutrally buoyant particles
- Bubbles are a poor choice: always overrespond unless quite small

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